

Electroweak Physics

Lecture II:

Status of Electroweak Physics

Acknowledgement:

Slides from M. Gruenewald, P. Renton, F. Teubert

Review

- Introduction to electromagnetic and weak interactions
- Motivation for Electroweak Unification
- Introduced nomenclature for electroweak studies
- Described electron-positron collisions and implications of the data

Helicity Conservation

Extreme Relativistic Limit (ERL): $E \gg mc^2$ $E = pc$ $\gamma \gg 1$

Massless limit

Helicity = chirality

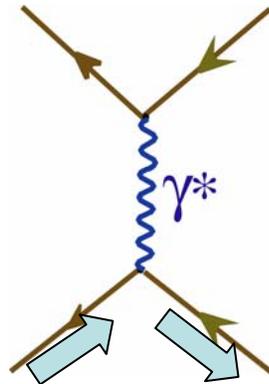
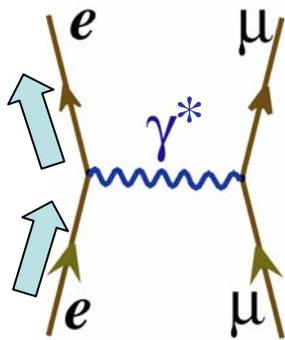
$$P_L \equiv \frac{(1 - \gamma^5)}{2} \quad P_R \equiv \frac{(1 + \gamma^5)}{2} \quad P_i P_j = \delta_{ij} P_j \quad \sum_i P_i = I \quad P_{L,R} u \equiv u_{L,R}$$

$$J_\mu^{EM} = q \bar{u} \gamma_\mu u = q (\bar{u}_L + \bar{u}_R) \gamma_\mu (u_L + u_R)$$

But $\bar{u}_L \gamma_\mu u_R = \bar{u}_R \gamma_\mu u_L = 0$

$$J_\mu^{EM} = q \bar{u}_R \gamma_\mu u_R + q \bar{u}_L \gamma_\mu u_L$$

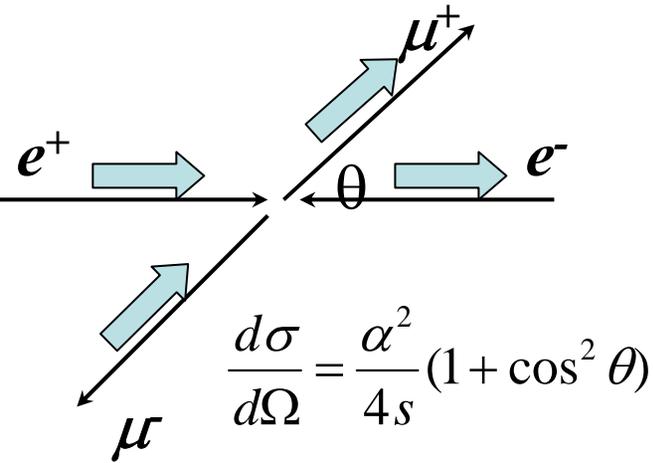
What are the implications?



For particle-antiparticle collisions

$$e^-_L + e^+_R \text{ or } e^-_R + e^+_L$$

Angular Distribution

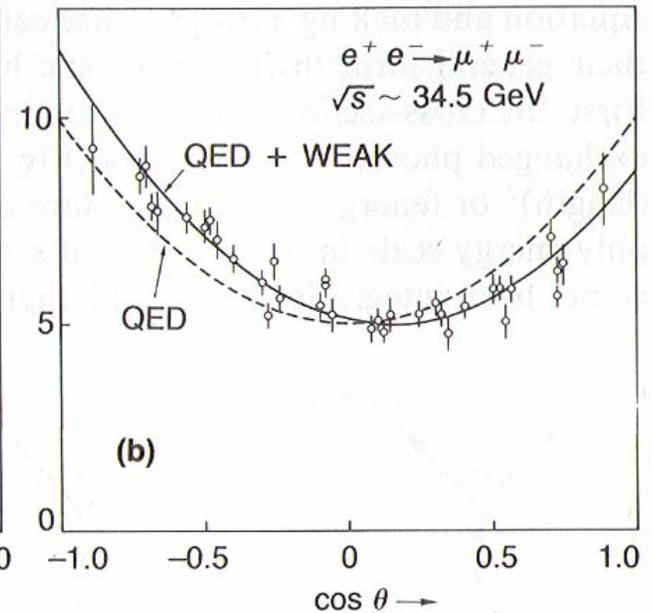
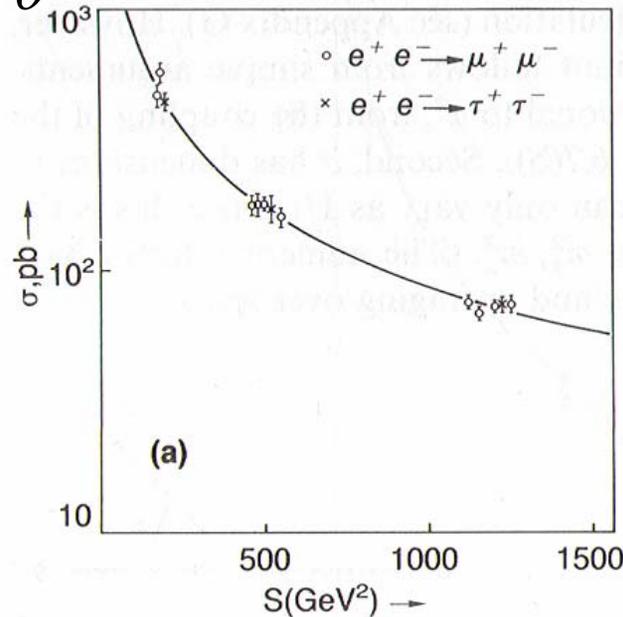


Start with spin 1 (forward or backward) along axis of collision: what is the probability of getting +1 or -1 along θ ?

$$d_{\lambda\lambda'}^j(\theta) \equiv \langle j\lambda' | e^{-i\theta J_y} | j\lambda \rangle$$

$$d_{11}^1(\theta) = d_{-1-1}^1(\theta) = \frac{1}{2}(1 + \cos \theta) \quad d_{1-1}^1(\theta) = d_{-11}^1(\theta) = \frac{1}{2}(1 - \cos \theta)$$

$$\sum (d_{\lambda\lambda'}^j(\theta))^2 = 1 + \cos^2 \theta$$



Z Decays

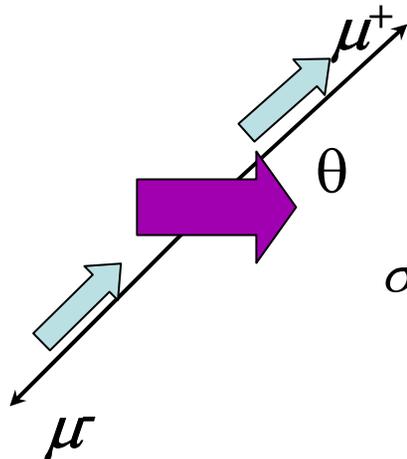
$$e^+e^- \rightarrow Z^0 \rightarrow l^+l^-, q\bar{q}$$

$$J_\mu^Z \sim g_R \bar{u}_R(e) \gamma_\mu u_R(e) + g_L \bar{u}_L \gamma_\mu u_L$$



Even if electrons and positrons are unpolarized, the Z's are produced polarized

$$P_Z = \frac{N_+ - N_-}{N_+ + N_-} = \frac{g_R^2 - g_L^2}{g_R^2 + g_L^2}$$



A_{++} : probability of $J = +1$ Z boson producing a $J = +1$ final state

$$A_{++} = \frac{g_R^e g_R^\mu}{2} (1 + \cos \theta)$$

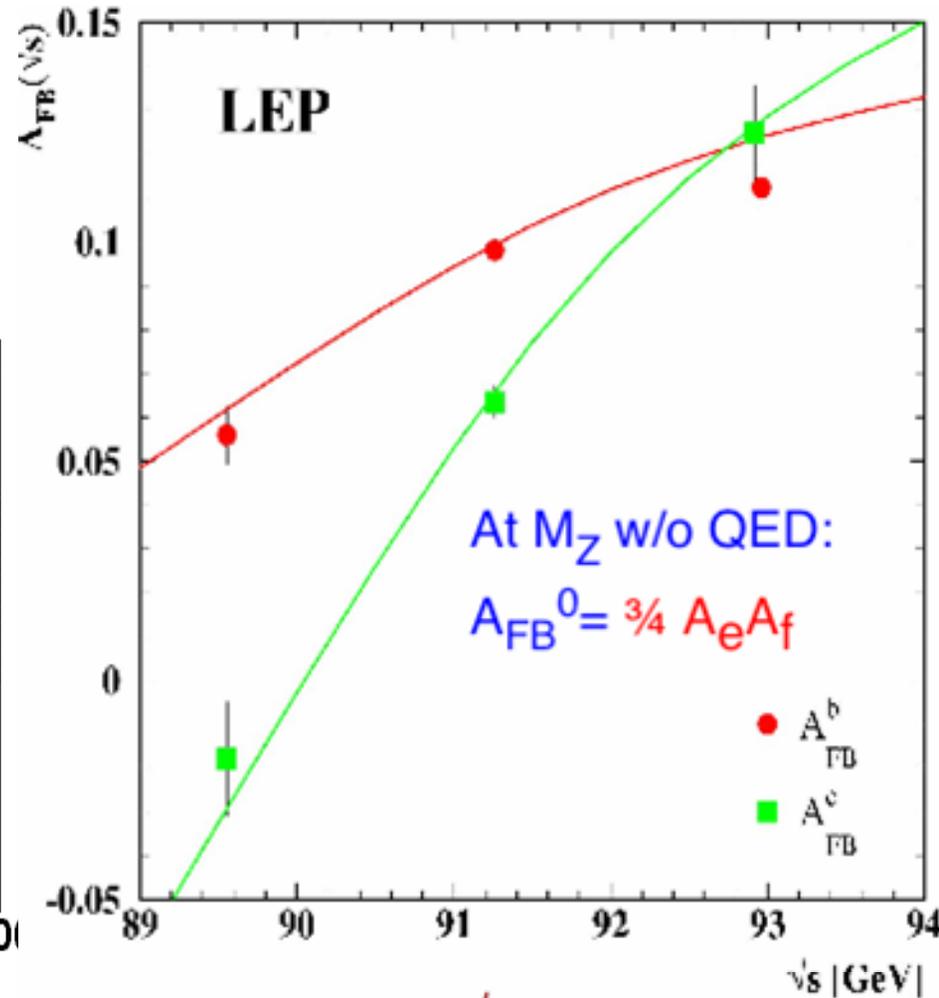
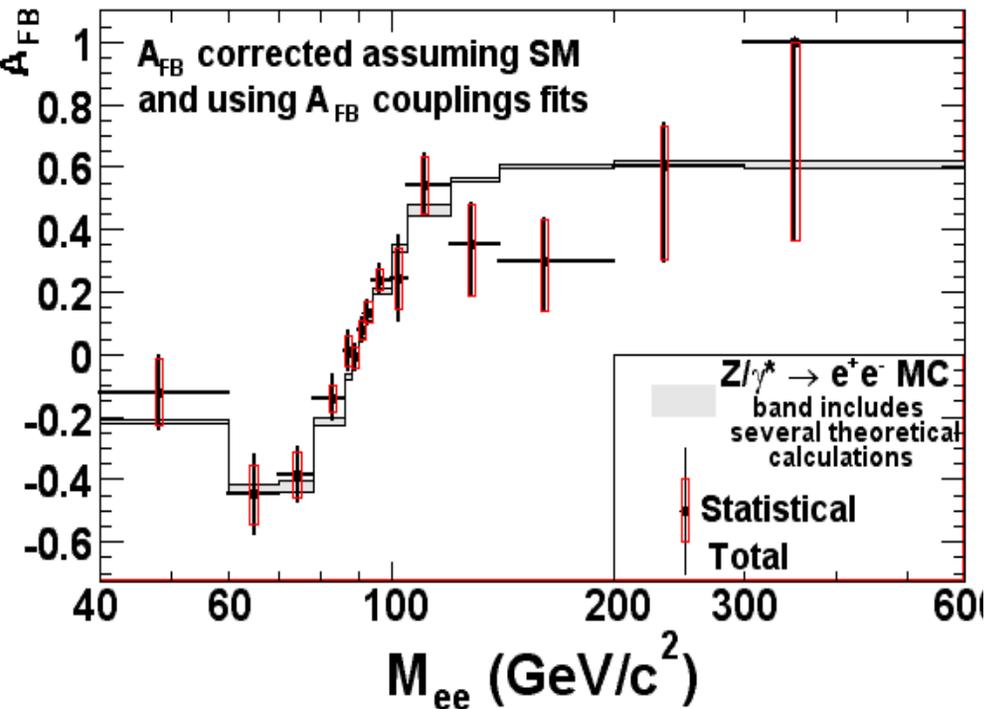
$$\sigma \propto \sum A_{ij}^2 = [(g_R^e)^2 + (g_L^e)^2][(g_R^\mu)^2 + (g_L^\mu)^2](1 + \cos^2 \theta) + [(g_R^e)^2 - (g_L^e)^2][(g_R^\mu)^2 - (g_L^\mu)^2] 2 \cos \theta$$

$$\sigma \propto (1 + \cos^2 \theta) + 2P_e P_\mu \cos \theta \quad P_e = \frac{(g_R^e)^2 - (g_L^e)^2}{(g_R^e)^2 + (g_L^e)^2} \quad P_\mu = \frac{(g_R^\mu)^2 - (g_L^\mu)^2}{(g_R^\mu)^2 + (g_L^\mu)^2}$$

Forward Backward Asymmetry

$$P_f = \frac{2g_{Vf}g_{Af}}{g_{Vf}^2 + g_{Af}^2} \approx 2 \frac{g_{Vf}}{g_{Af}} \approx 1 - 4 \sin^2 \theta_w$$

$$A_{FB} = \frac{\sigma(\cos \theta > 0) - \sigma(\cos \theta < 0)}{\sigma(\cos \theta > 0) + \sigma(\cos \theta < 0)} = \frac{3}{4} P_e P_f$$



$$A_f = 2 \frac{g_{Vf} / g_{Af}}{1 + (g_{Vf} / g_{Af})^2} \Leftrightarrow \sin^2 \theta_{eff}$$

Left-Right Asymmetry

AFB is the product of 2 small numbers.

Can they be disentangled?

Polarize the electron beam and measure Z production

$$P_b = \frac{N_+ - N_-}{N_+ + N_-} \quad \text{Fraction of beam polarized along or against the momentum}$$

$$A_{LR} = \frac{N_{Z^-} - N_{Z^+}}{N_{Z^-} + N_{Z^+}} = \frac{(1 - P_b)g_L^2 - (1 + P_b)g_R^2}{(1 - P_b)g_L^2 + (1 + P_b)g_R^2} = P_b P_e$$

All final states can be used!

This was the motivation for the SLAC Linear Collider:

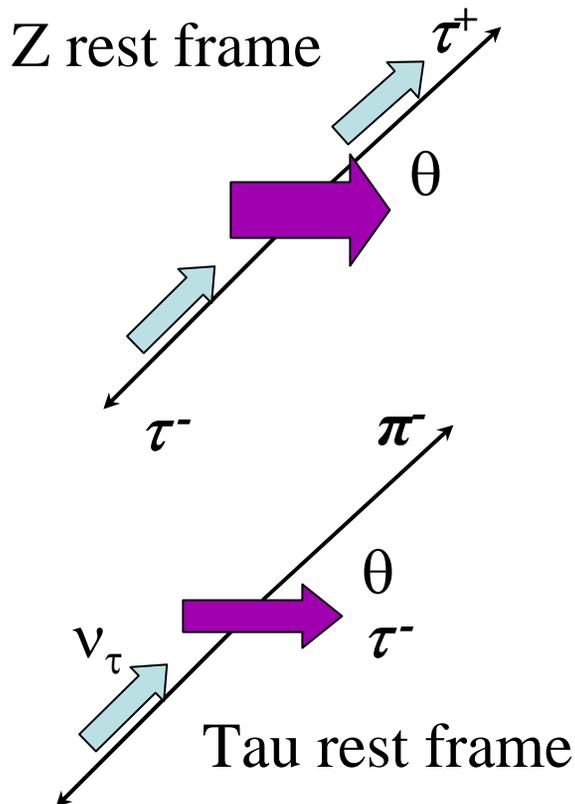
Could compete with a factor of 10 to 100 less luminosity

Tau Polarization

Instead of polarizing the initial state,

Analyze the final state: need a polarization filter

For tau leptons, use the weak decay!



$$\Gamma_\tau = \frac{G_F^2 m_\tau^5}{192\pi}$$

Lifetime ~ few ps

Travels a few mm

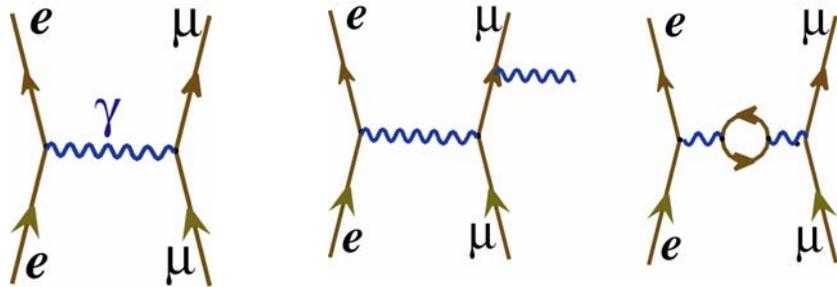
V-A interaction reveals tau polarization

Pion lab energy distribution is related trivially to the rest frame angular distribution

Perturbation Theory

From Feynman rules: Construct all possible diagrams

Consistent with standard conservation laws

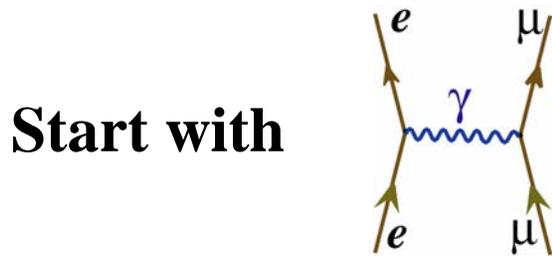


Amplitude is sum of all possible states: Feynman's path integral formulation of QM

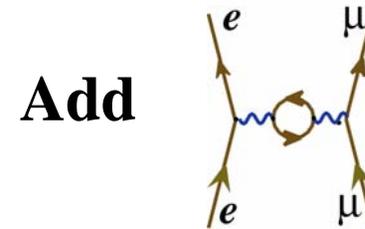
Problem: Total amplitude diverges

- *Feynman rules with electric charge*
- *Calculate $\sigma_1(e)$ for a test process*
- *Measure $\sigma_1(e)$ and extract e*
- *Calculate $\sigma_2(e)$ for another process*

Charge Renormalization



$$M_1 \sim \frac{e^2}{q^2}$$

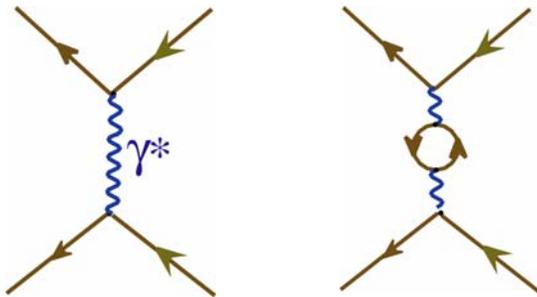


Introduce parameter $\Sigma_{\gamma\gamma}(q^2)$

(It is infinite)

$$M_1 \sim i \frac{e^2}{q^2} + i \frac{e^2}{q^2} \frac{i \Sigma_{\gamma\gamma}(q^2)}{q^2}$$

$$e^2 \longrightarrow e^2 (1 - \Pi_{\gamma\gamma}(q^2))$$



Introduce parameter $\Sigma_{\gamma\gamma}(p^2)$ **(Also infinite)**

$$M_2 \sim \frac{e^2}{p^2} (1 - \Pi_{\gamma\gamma}(p^2))$$

$$M_2 \sim \frac{e^2}{p^2} (1 - [\Pi_{\gamma\gamma}(p^2) - \Pi_{\gamma\gamma}(q^2)]) \quad \textbf{Finite!}$$

Running Couplings

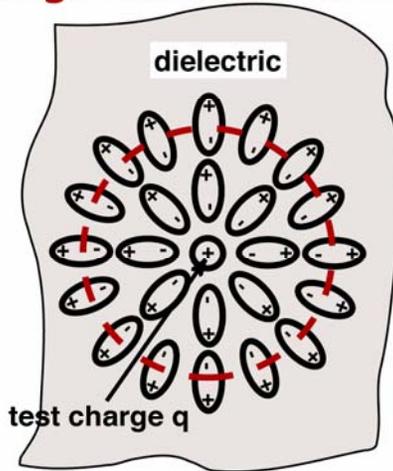
Fine structure constant: $1/137$ at low energy, $1/128$ at Z pole

Not all Quantum Field Theories behave this way:
The ones that do are renormalizable theories

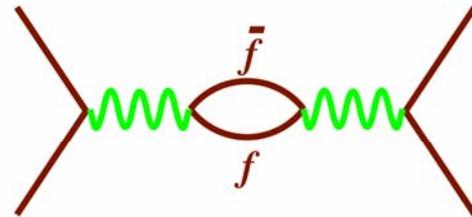
Electroweak theory: t'Hooft and Veltman

QCD: Gross, Politzer and Wilzcek

total charge enclosed is less than q



total charge depends on relative distance

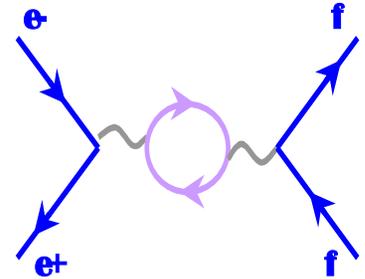


*effective charge increases
with decreasing distance:*

**higher order terms in
perturbative expansion**

Calculation of Running

The shift $\Delta\alpha$ can be determined **analytically for lepton loops** and by **a dispersion integral over the e^+e^- annihilation cross section for light quarks** (u,d,s,c,b)



$$\Delta\alpha_{\text{lepton}} = \sum_{l=e,\mu,\tau} \frac{\alpha}{3\pi} \left(\log \frac{m_Z^2}{m_l^2} - \frac{5}{3} \right) + \dots$$

$$\alpha(m_Z^2) = \alpha / (1 - \Delta\alpha)$$

Optical theorem

$$\Delta\alpha_{\text{hadron}} = - \frac{\alpha}{3\pi} \int_{4m_\pi^2}^{\infty} \frac{m_Z^2 ds'}{s' [s' - m_Z^2]} \frac{\sigma(e^+e^- \rightarrow \gamma^* \rightarrow q\bar{q})}{\sigma(e^+e^- \rightarrow \gamma^* \rightarrow \mu^+\mu^-)}$$

Electroweak Input Parameters

For electroweak interactions, there are three parameters needed:

1. Scale of electromagnetism (electric charge)
2. Scale of the weak interaction (Vector boson mass)
3. Weak mixing angle

Parameters are chosen from experimental measurements:

1. Low energy Thomson Scattering
2. The muon lifetime
3. The mass of the Z boson

Z mass known to 23 parts per million!

LEP at CERN

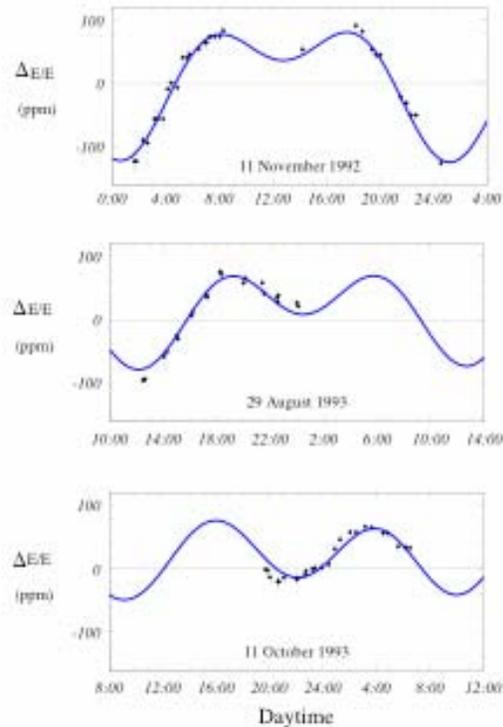


June 2, 2005

Electroweak Physics: Lecture II

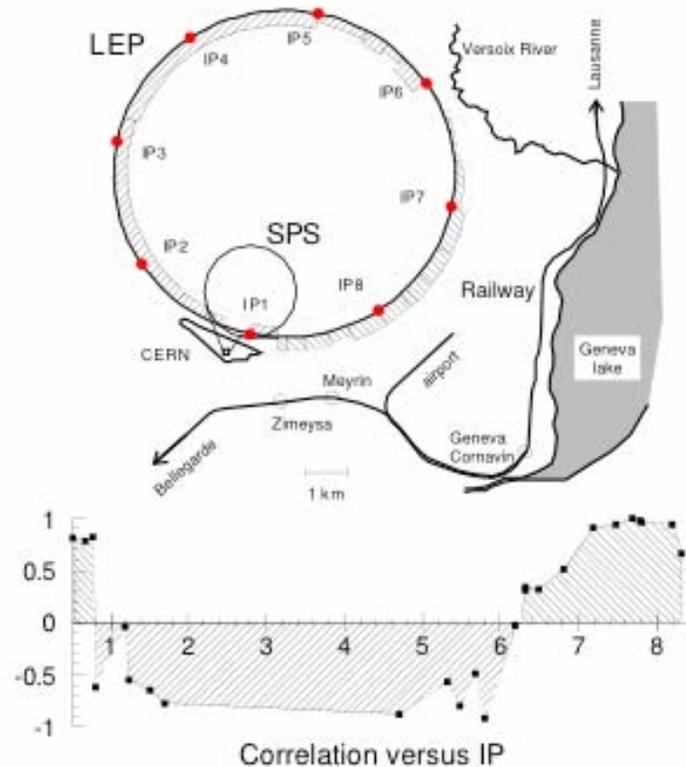
Main energy systematics

Tides and other ground motion:



Use tidal model and beam position monitors to correct for orbit changes

Trains

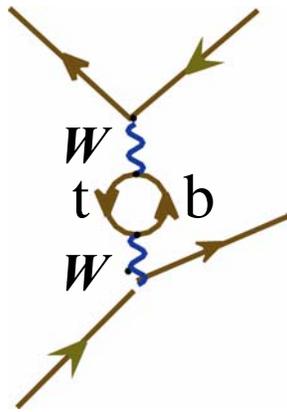


Use NMR probes and *thermal* model to extrapolate energy during fills

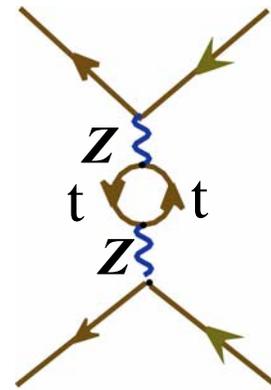
Indirect Evidence for the Top

Measure any asymmetry on the Z pole: function of weak mixing angle

The answer differs from what you would get at tree level



Muon decay



Z production

$$\Pi_{WW} - \Pi_{ZZ} \propto m_t^2 - m_b^2$$

Electroweak Precision Data

Very high Q^2 physics at LEP, SLC, and the Tevatron:

More than 1000 measurements with (correlated) uncertainties
Combined to 17 precision electroweak observables

Z boson physics (LEP-1, SLD):

- 5 Z lineshape and leptonic forward-backward asymmetries
- 2 Polarised leptonic asymmetries P_τ , $A_{LR}(FB)$
- 1 Inclusive hadronic charge asymmetry
- 6 Heavy quark flavour results (Z decays to b and c quarks)

W boson & top quark physics – ongoing at Tevatron's Run-II:

- 2 W boson mass and width (LEP-2, Tevatron)
- 1 Top quark mass (Tevatron)

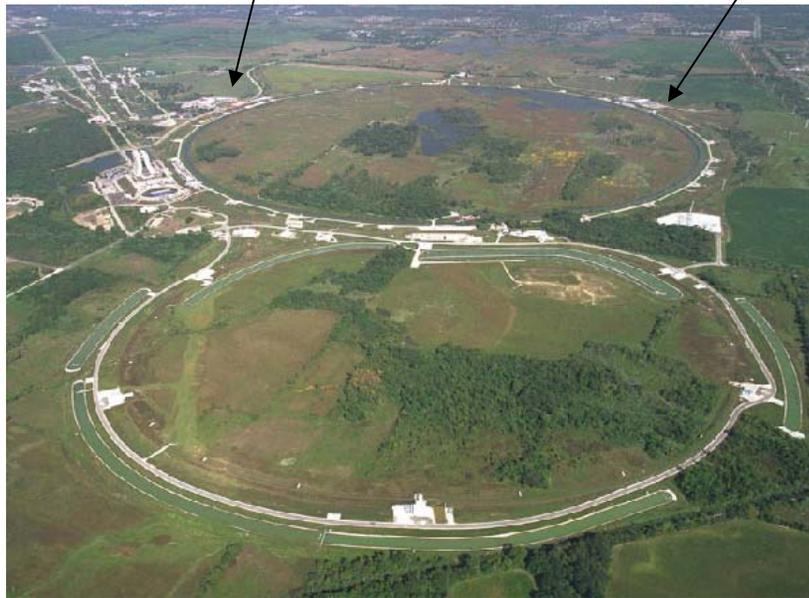
Fermilab



CDF



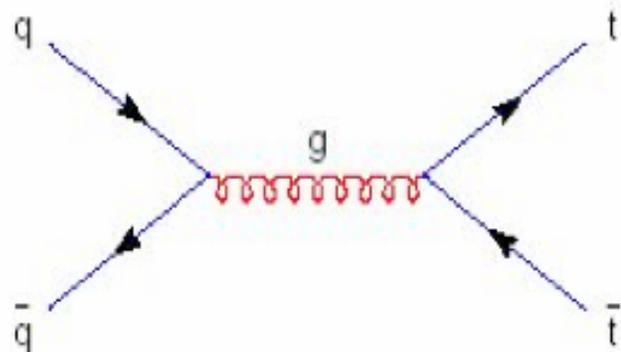
D0



Top Physics

Tevatron: only source of top quarks in the world!

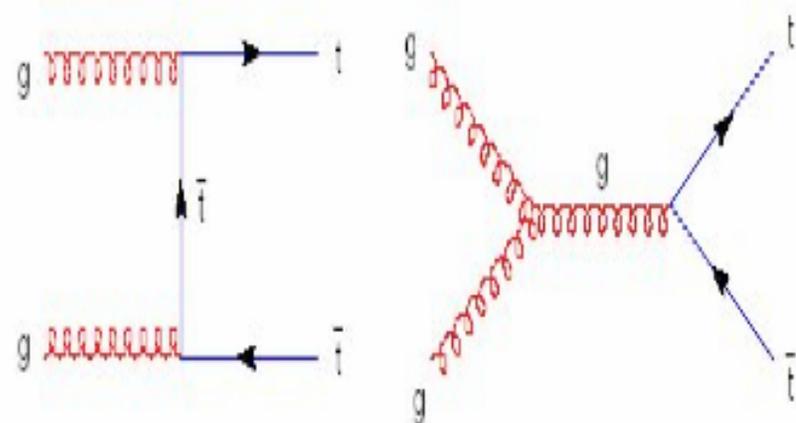
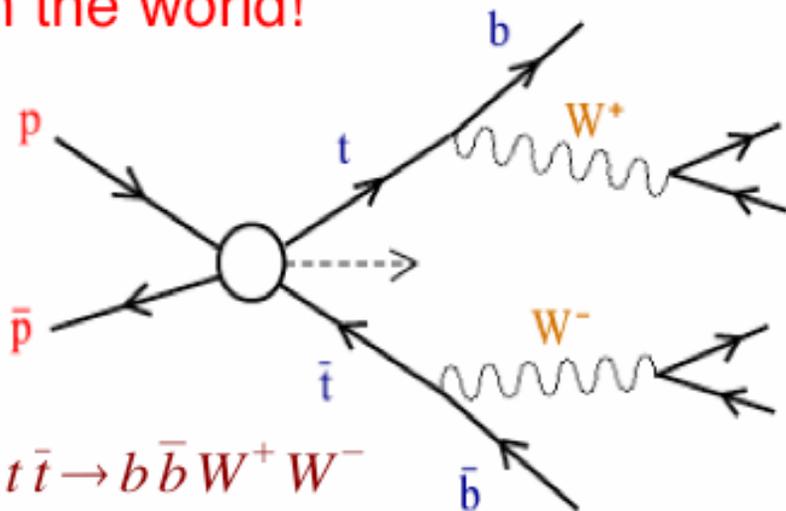
Primarily top-pair production



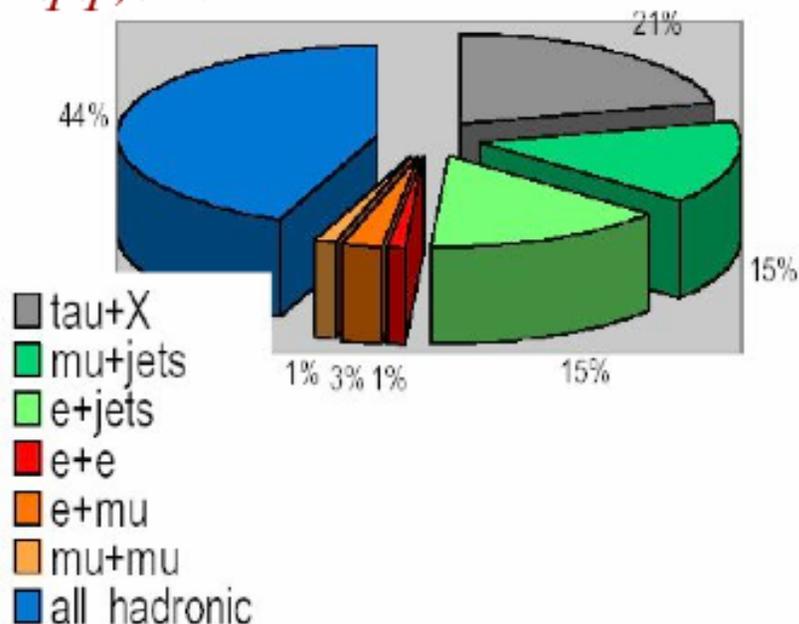
$q\bar{q}$ annihilation (85%)

$$p\bar{p} \rightarrow t\bar{t}X, \quad t\bar{t} \rightarrow b\bar{b}W^+W^-$$

$$W^- \rightarrow q\bar{q}, l^-\bar{\nu}$$

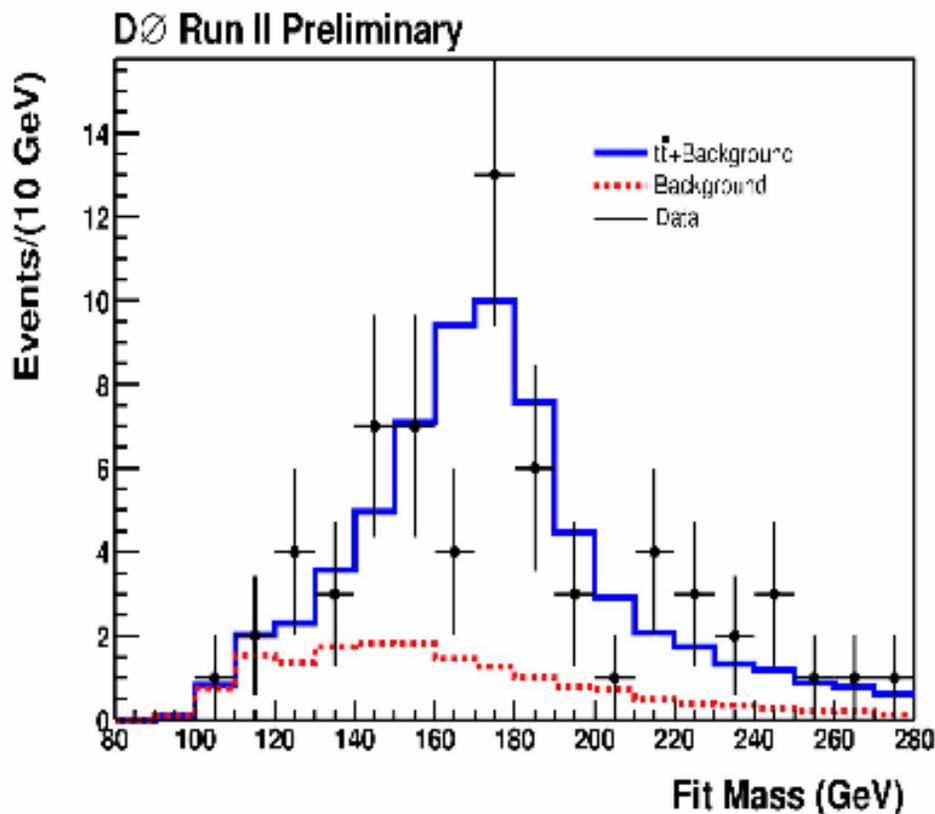
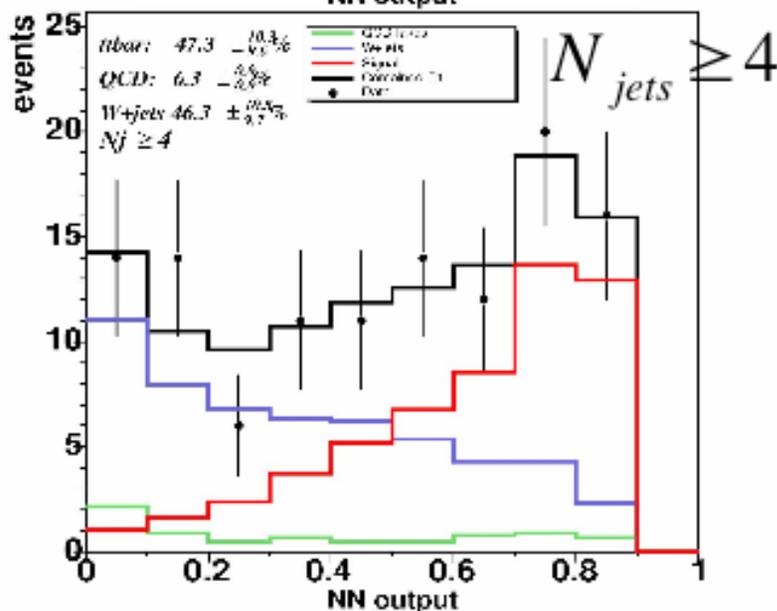
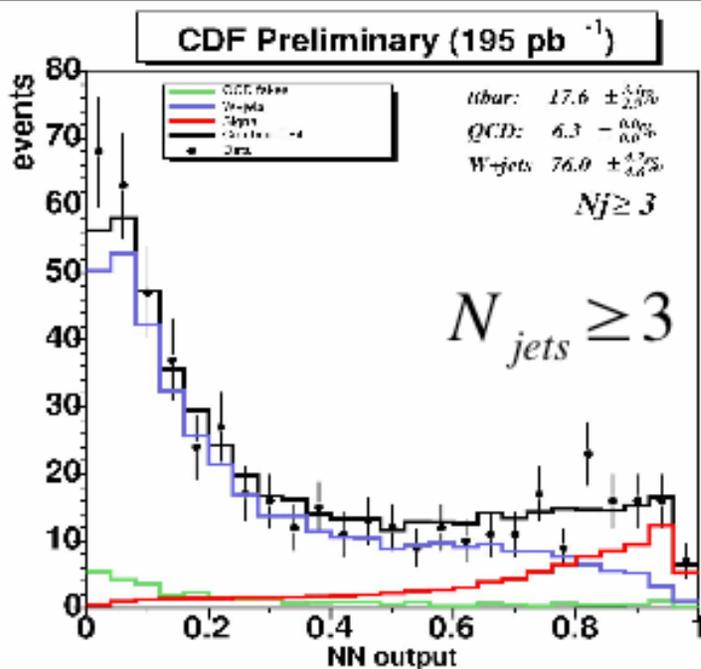


gluon fusion (15%)



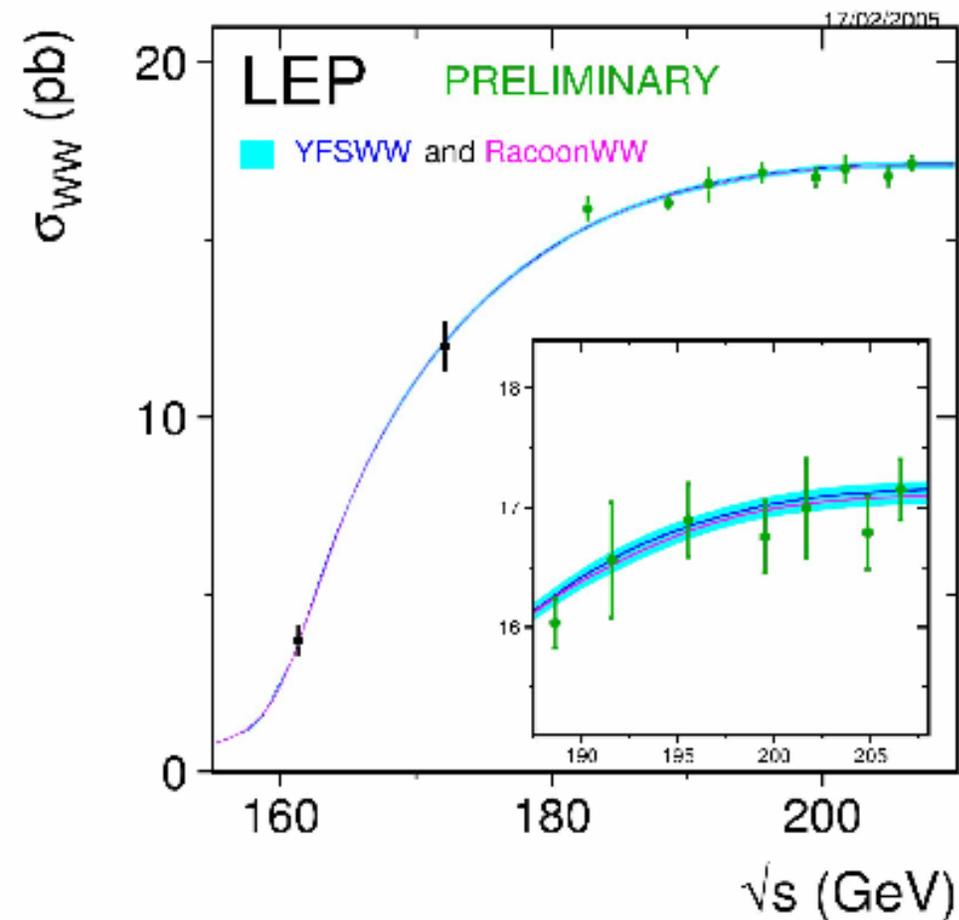
Top Production

Lepton+jets most promising channel:
 Charged lepton, 2 b-quark jets
 2 other jets, only 1 neutrino
 Invariant mass $M(\text{top}) = M(Wb)$



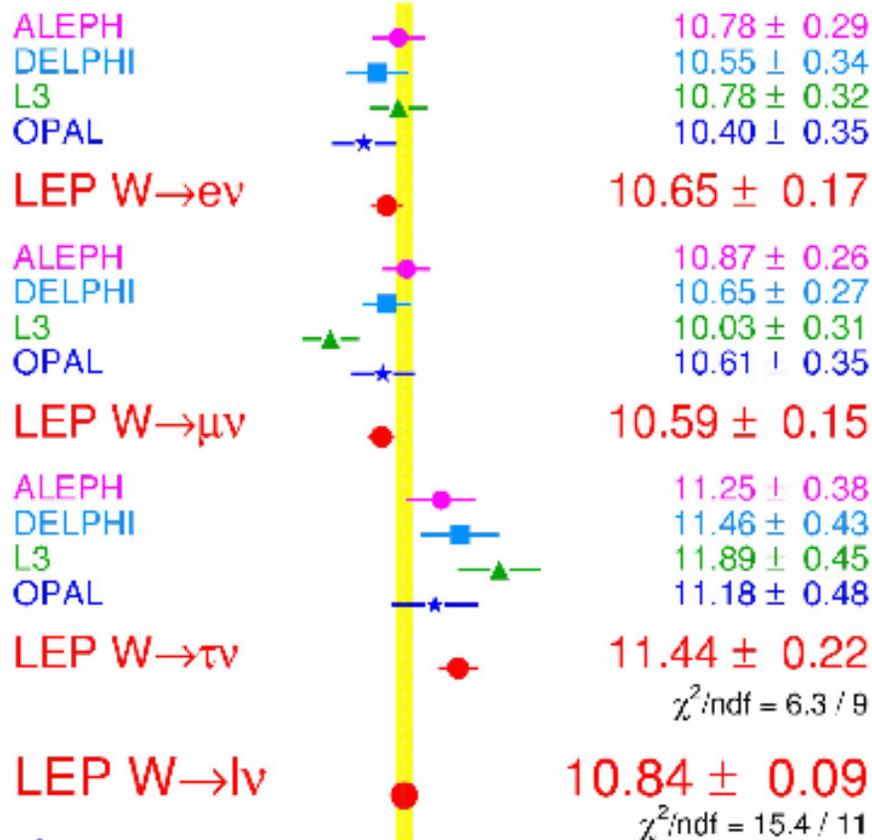
W-Pairs at LEP

Winter 2005 - LEP Preliminary



W Leptonic Branching Ratios

23/02/2005



ALEPH, DELPHI, L3 final, OPAL prel.

Subsequent maximisation of discrepancy:

W-tau branching fraction $\sim 2.9\sigma$ above W-e/ μ average

Standard Model Analysis

SM: Each observable calculated as a function of:

$\Delta\alpha_{\text{had}}, \alpha_s(M_Z), M_Z, M_{\text{top}}, M_{\text{Higgs}}$ (and G_F)

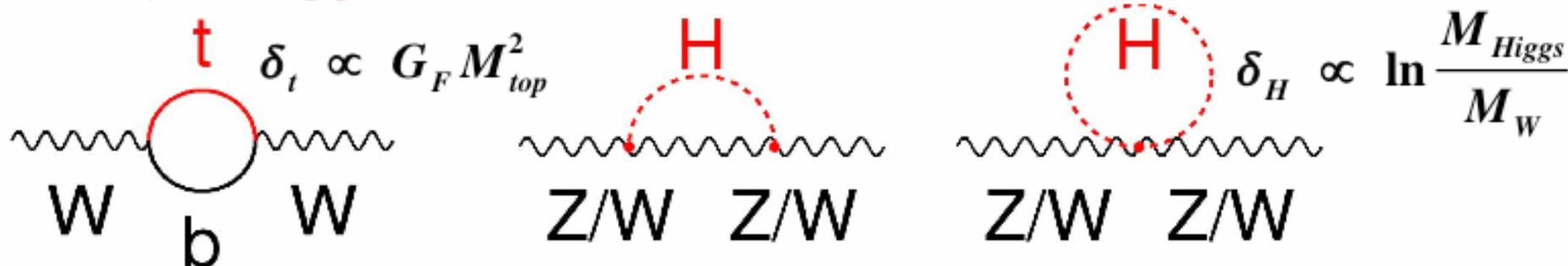
$\Delta\alpha_{\text{had}}$: hadronic vacuum polarisation $[0.02761 \pm 0.00036]$

$\alpha_s(M_Z)$: given by Γ_{had} and related observables

M_Z : constrained by LEP-1 lineshape

Precision requires 1st and 2nd order electroweak and mixed radiative correction calculations (QED to 3rd)

$M_{\text{top}}, M_{\text{Higgs}}$ enter through electroweak corrections ($\sim 1\%$)!



Calculations by programs TOPAZ0 and ZFITTER

Heavy Particle Masses W and Top

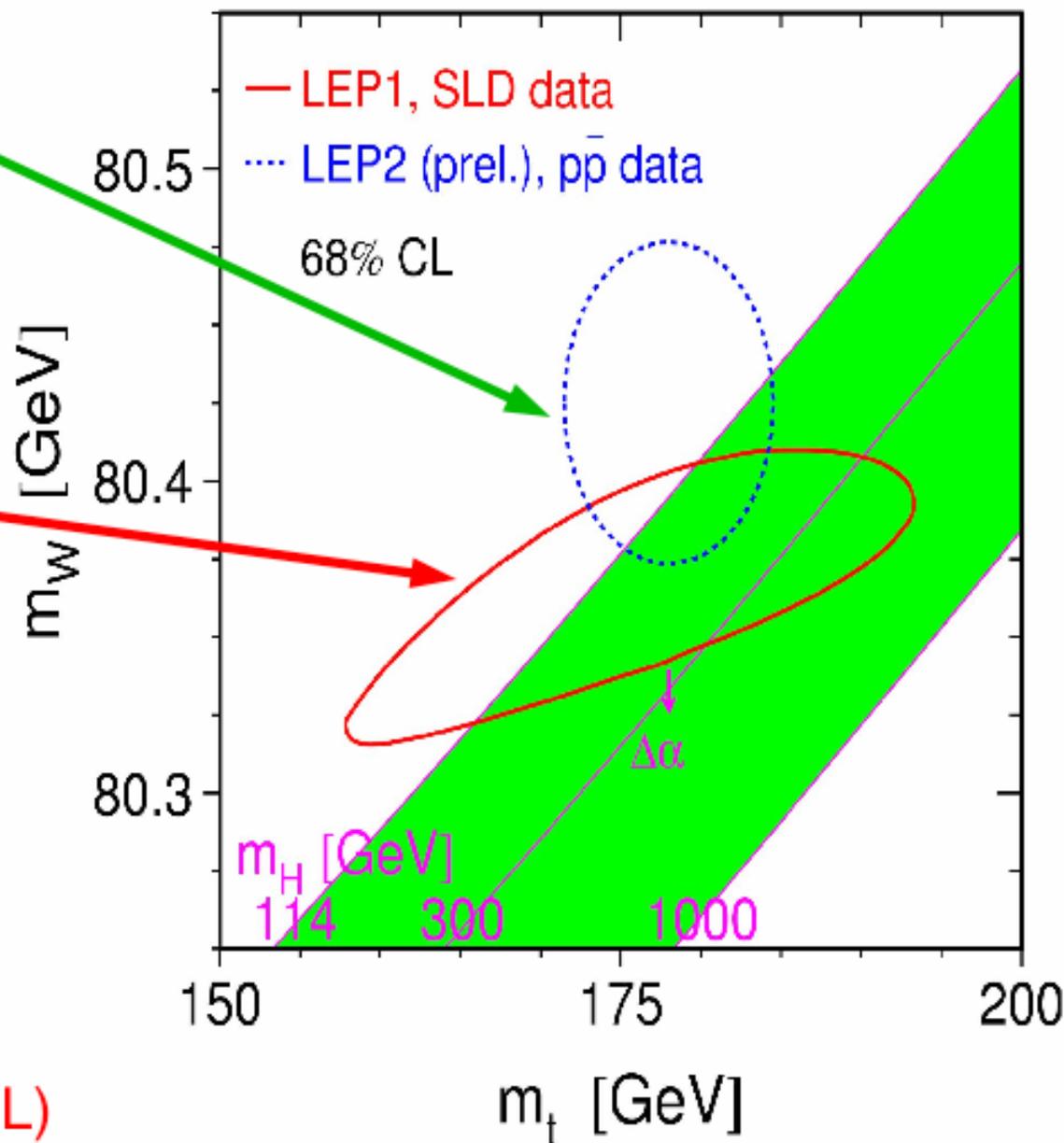
Direct measurements:
Tevatron and LEP2

Z-Pole measurements:
Constrain electroweak
radiative corrections
Allow to predict M_W
and M_{top} within SM

Good agreement:
Successful SM test

Both data sets prefer a
light Higgs boson

$M_{\text{Higgs}} < 280 \text{ GeV}$ (95%CL)



Standard Model Analysis

$$M_{\text{Higgs}} = 126^{+73}_{-48} \text{ GeV}$$

Incl. theory uncertainty:

$$M_{\text{Higgs}} < 280 \text{ GeV (95\%CL)}$$

does not include:

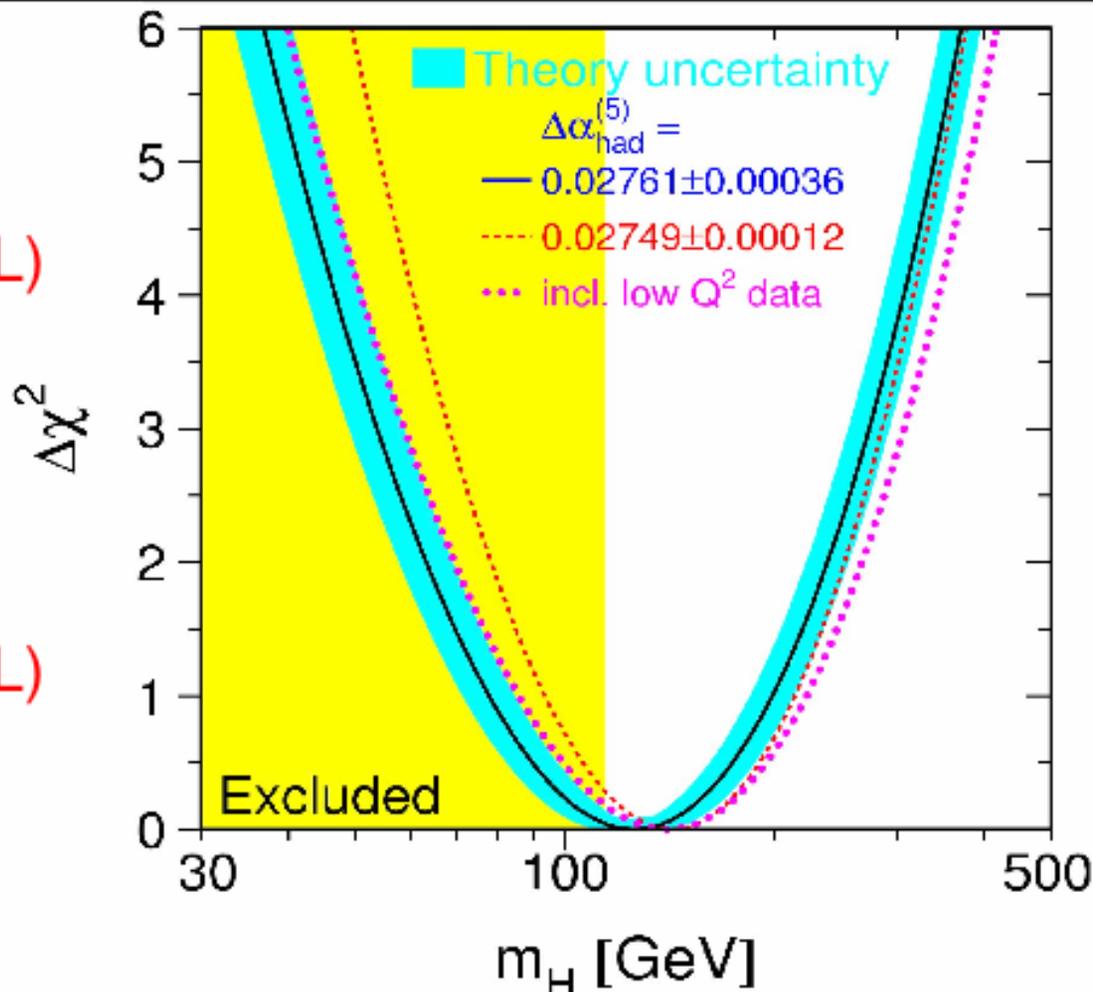
Direct search limit (LEP-2):

$$M_{\text{Higgs}} > 114 \text{ GeV (95\%CL)}$$

Renormalise probability

for $M_{\text{H}} > 114 \text{ GeV}$ to 100%:

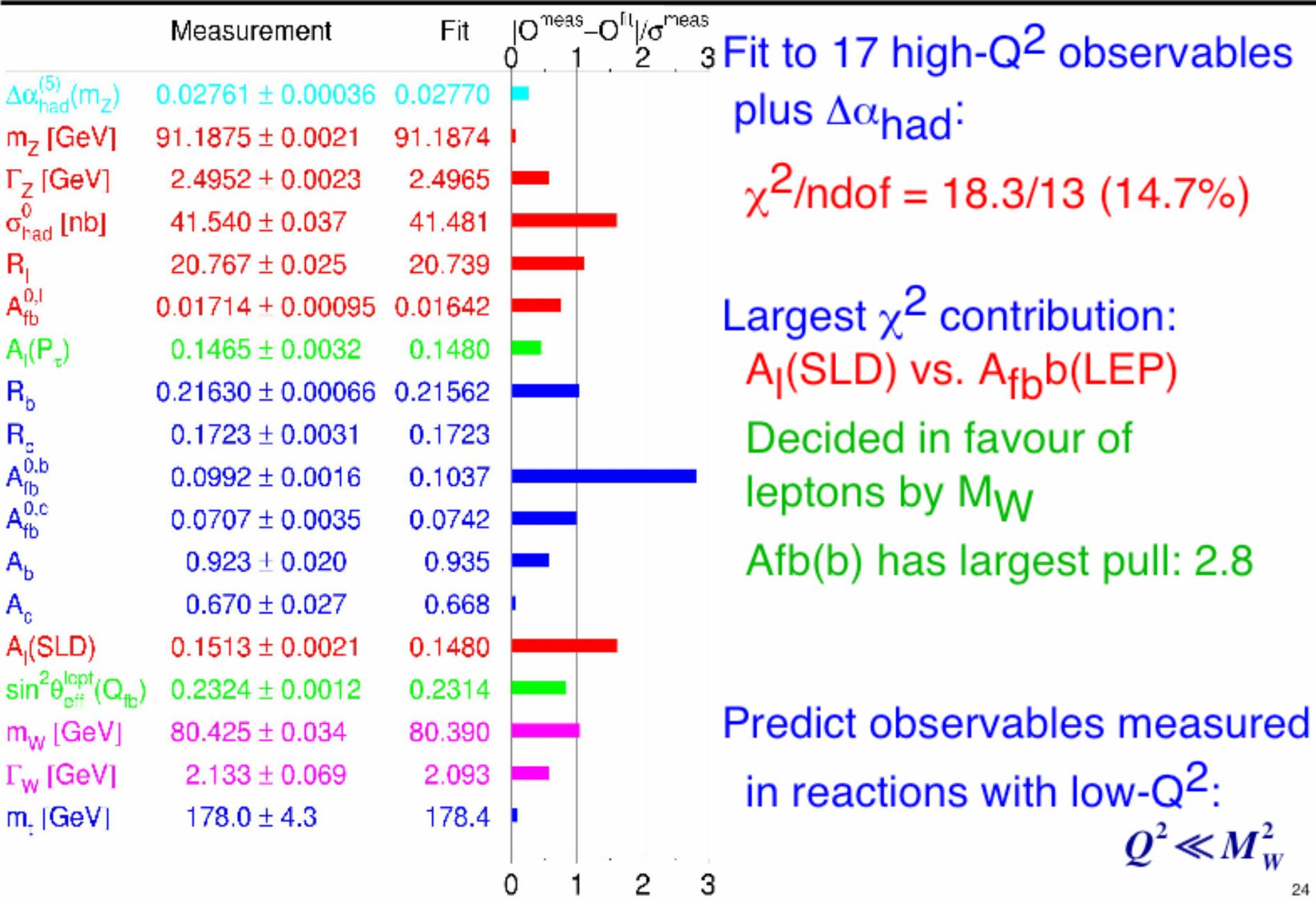
$$M_{\text{Higgs}} < 300 \text{ GeV (95\%CL)}$$



Theory uncertainty:

Dominated by two-loop
calculations for $\sin^2\Theta_{\text{eff}}$

Standard Model Analysis



Comparison of all Z-Pole Asymmetries

Effective electroweak
mixing angle:

$$\sin^2\theta_{\text{eff}} = (1 - g_V/g_A)/4$$

$$= 0.23153 \pm 0.00016$$

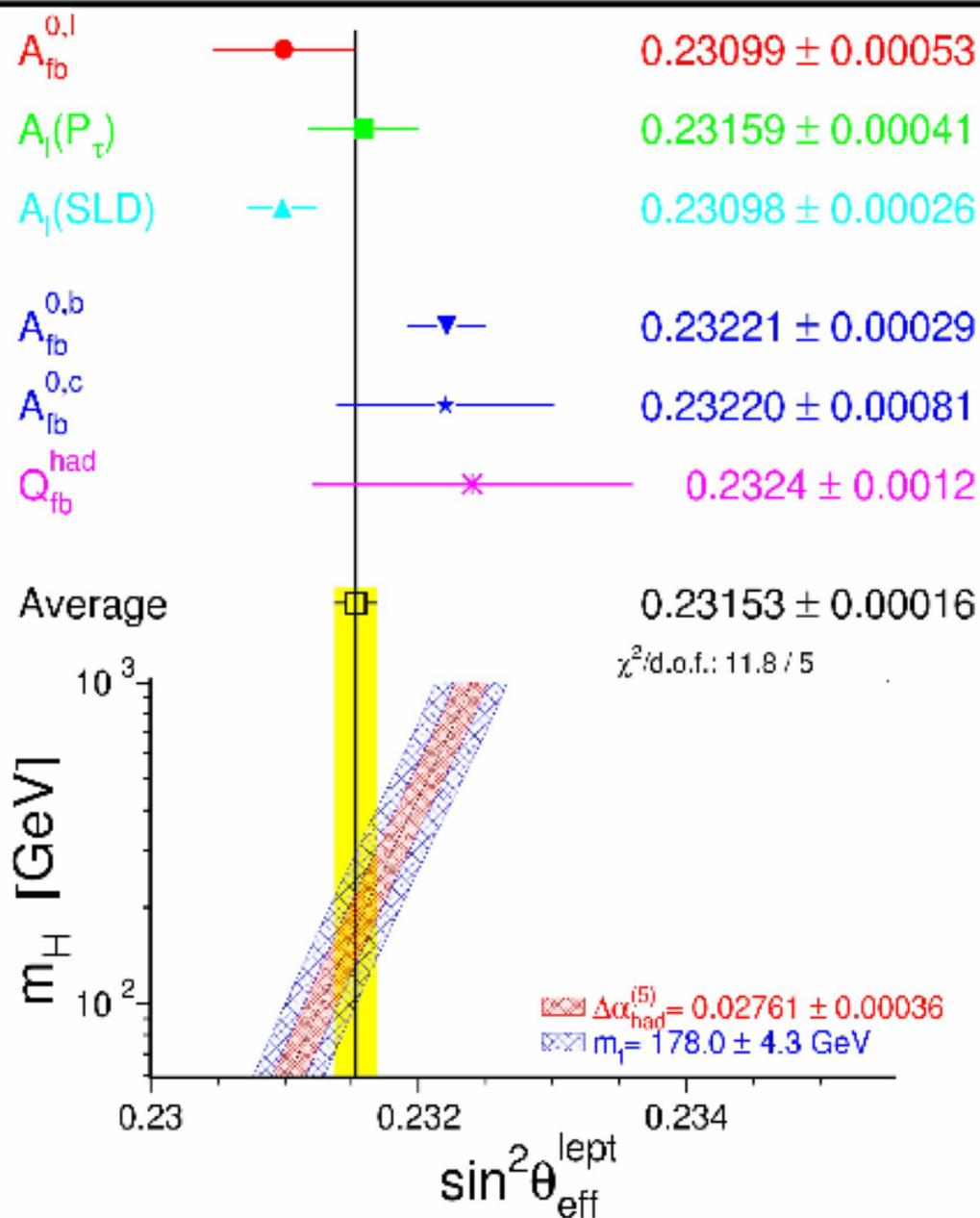
$$\chi^2/\text{ndof} = 11.8/5 \quad [3.8\%]$$

Subsequent observation:

0.23113 ± 0.00021 leptons
0.23222 ± 0.00027 hadrons
3.2 σ difference

But is really:

$A_1(\text{SLD})$ vs. $A_{\text{fb}}^b(\text{LEP})$
3.2 σ difference



Summary

- The electroweak theory has been tested to extraordinary precision
- Typical scale is 0.1%
- No significant deviations, but some tantalizing hints
- The Higgs boson is expected to be light
- Should we pack up and go home? Some answers in Lecture III