

JLab: Probing Hadronic Physics with Electrons and Photons

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Jefferson Lab

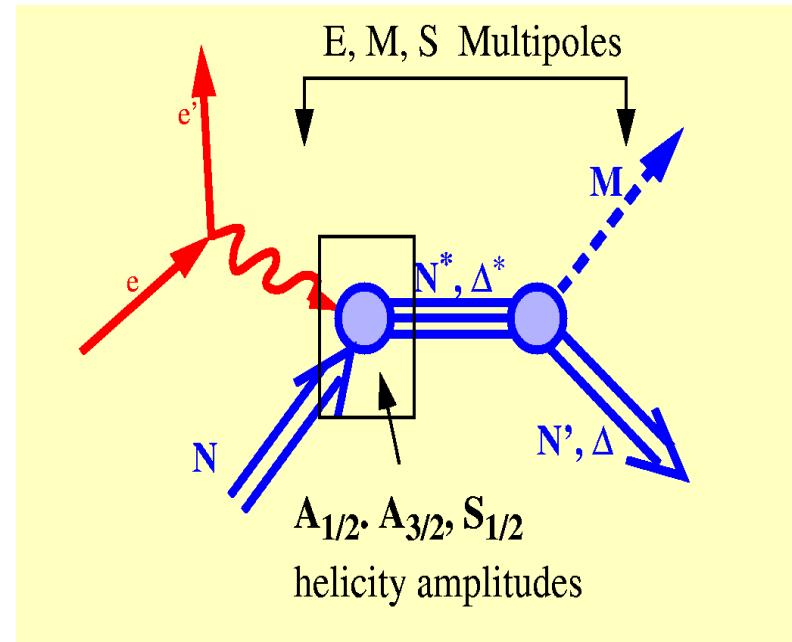
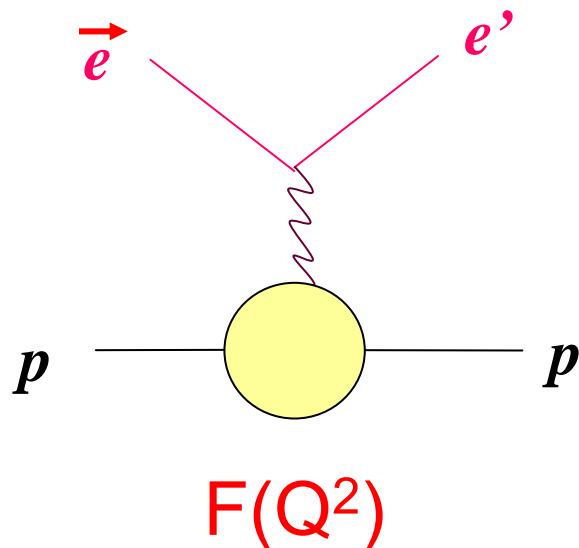
V Latinamerican Symposium on
Nuclear Physics

Santos, Brazil, September 2003

Introduction to JLab
The shape of the proton
Pentaquarks

Why use electron and photon probes?

Electromagnetic interaction is well-known



Elastic Form Factors

Inelastic transitions

Size probed $\sim 1/\sqrt{Q^2}$

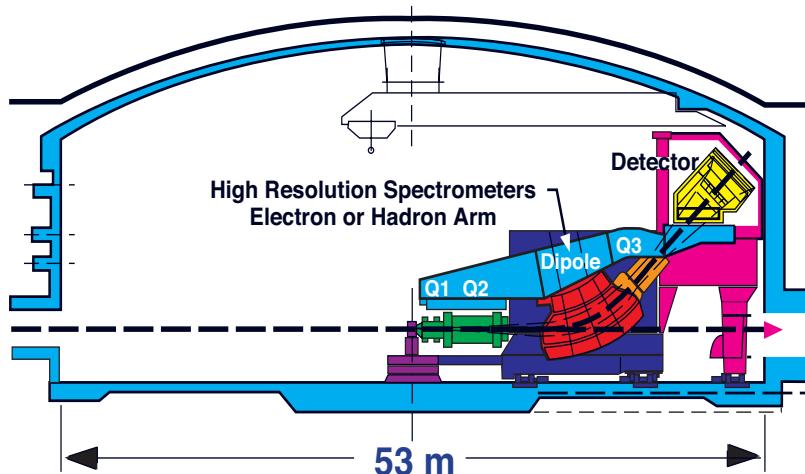
CEBAF @ JLab Today

- Main physics programs
 - nucleon electromagnetic form factors (including strange form factors)
 - $N \rightarrow N^*$ electromagnetic transition form factors
 - spin structure functions of the nucleon
 - form factors and structure of light nuclei
 - Superconducting recirculating electron accelerator
 - max. energy 5.7 GeV
 - max current 200 μ A
 - e polarization 80%
 - Simultaneous operation in 3 halls $L[\text{cm}^{-2}\text{s}^{-1}]$
 - A: Two High Resolution Spectrometers ($p_{\text{max}}=4$ GeV/c) 10^{39}
 - B: Large Acceptance Spectrometer for e and γ induced reactions 10^{34}
 - C: Two spectrometers ($p_{\text{max}}=7$ and 1.8 GeV/c) + special equipment 10^{39}

CEBAF accelerator site

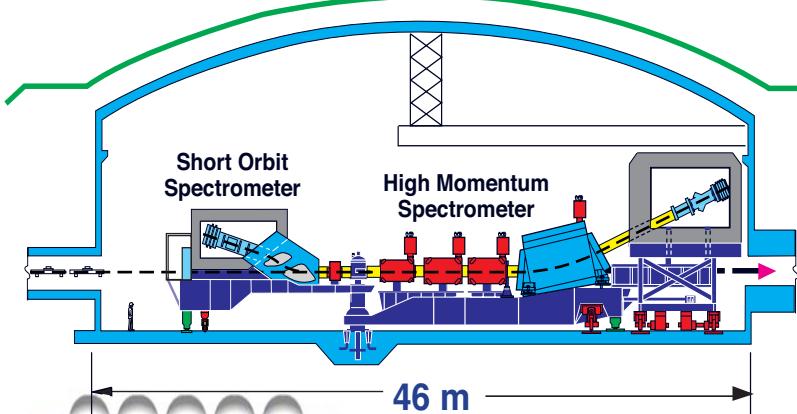


Three Experimental End-Stations



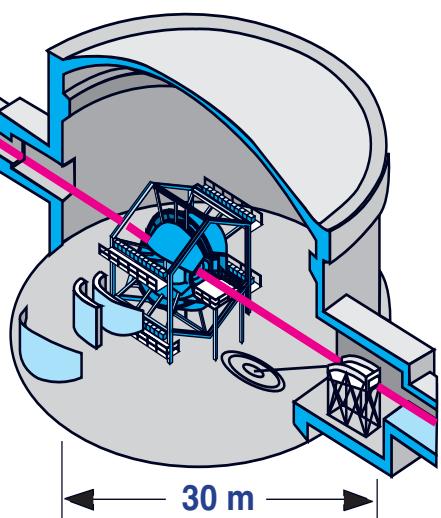
HALL A

Pair of identical High Resolution Spectrometers (HRS²)



HALL C

High Momentum Spectrometer (HMS) and Short Orbit Spectrometer (SOS)



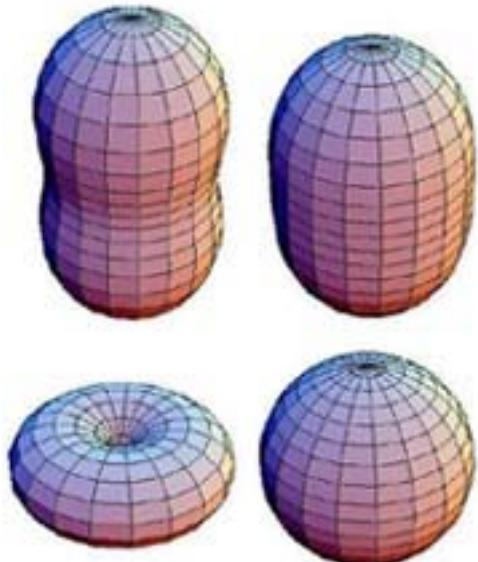
HALL B

CEBAF's Large Acceptance Spectrometer (CLAS) and Bremsstrahlung Photon Tagger

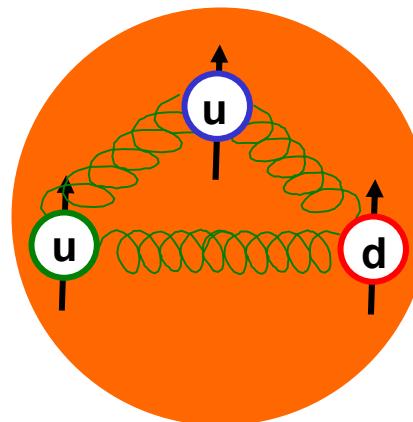
G_{Ep} : Electric form factor of the proton

Goal: Determine the charge and current distributions inside the proton

NY Times “Is a proton round?”



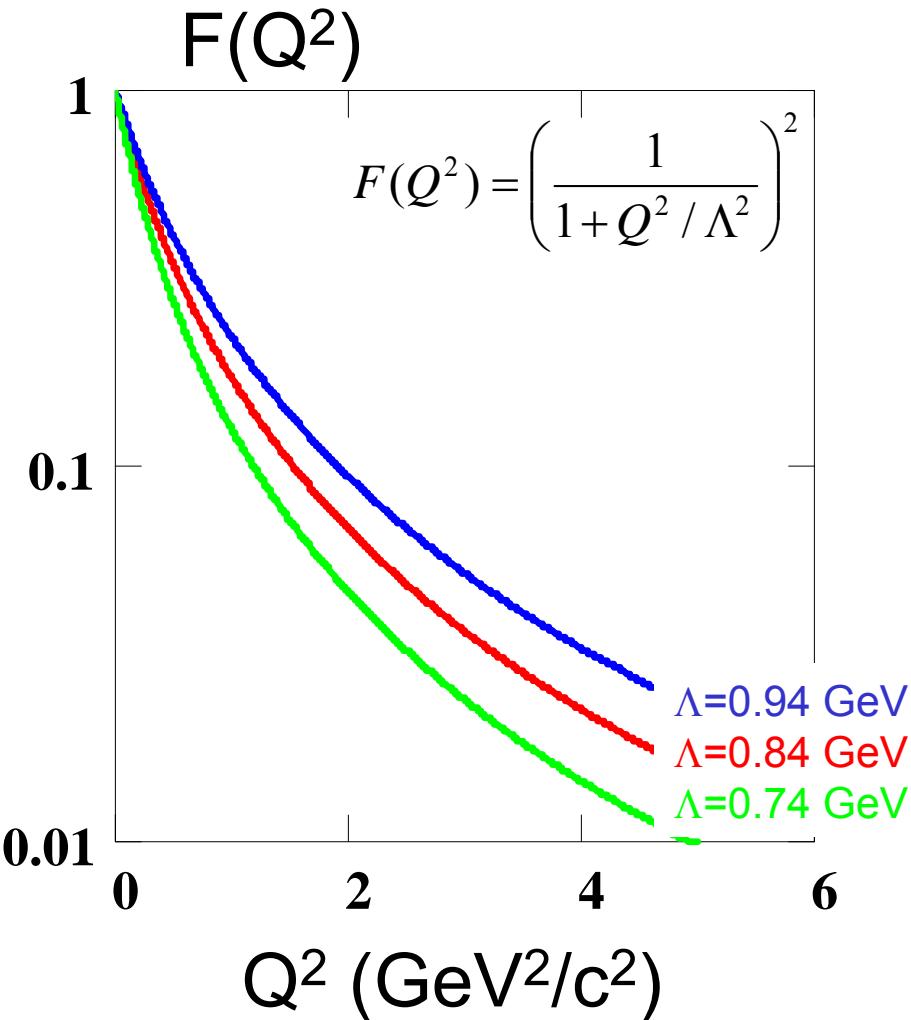
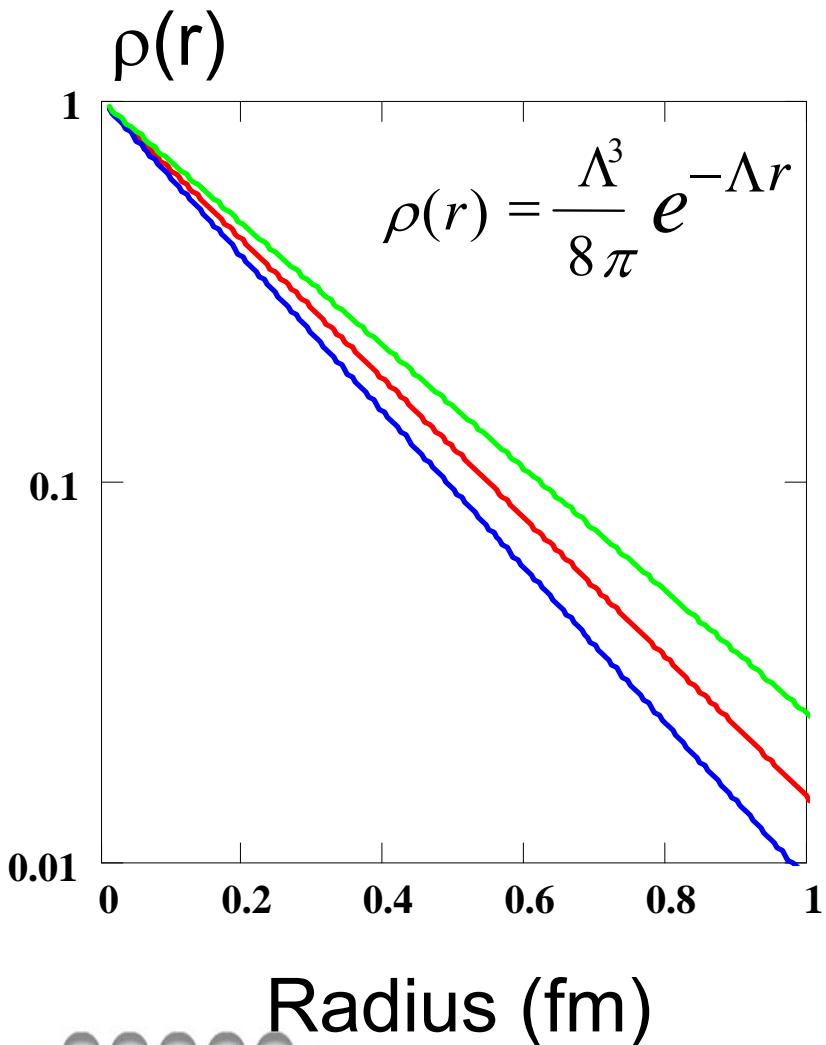
from G.A. Miller



Naïve expectation is that the charge and currents are determined from the spatial distribution of quark charges and spins.

Charge distribution and Form Factors

$$\rho(r) = \int d^3q F(q^2) e^{-i\vec{q}\cdot\vec{r}}$$



Charge Distributions and Form Factors

point

$$f(r) = \delta(r - r_o)$$

$$F(Q^2) = 1$$

unity

exponential

$$f(r) = \frac{\Lambda^3}{8\pi} e^{-\Lambda r}$$

$$F(Q^2) = \left(\frac{1}{1 + Q^2 / \Lambda^2} \right)^2$$

dipole

Yukawa

$$f(r) = \frac{\Lambda^2}{4\pi r} e^{-\Lambda r}$$

$$F(Q^2) = \frac{1}{1 + Q^2 / \Lambda^2}$$

pole

Gaussian

$$f(r) = \left(\frac{\Lambda^2}{2\pi} \right) e^{-\left(\frac{1}{2} \Lambda^2 r^2 \right)}$$

$$F(Q^2) = e^{-\frac{1}{2} \left(Q^2 / \Lambda^2 \right)}$$

Gaussian

Decomposition of the elastic cross section

$$\frac{d\sigma}{d\Omega} = \sigma_{ns} \left(\frac{\mathbf{G}_{Ep}^2 + \tau \mathbf{G}_{Mp}^2}{1 + \tau} + 2 \tau \mathbf{G}_{Mp}^2 \tan^2(\vartheta_e / 2) \right)$$

$$\sigma_{ns} = \frac{\alpha^2 \cos^2(\vartheta_e / 2)}{4E^2 \sin^4(\vartheta_e / 2)} \frac{E'}{E} \quad \tau = \frac{Q^2}{4M_p^2}$$

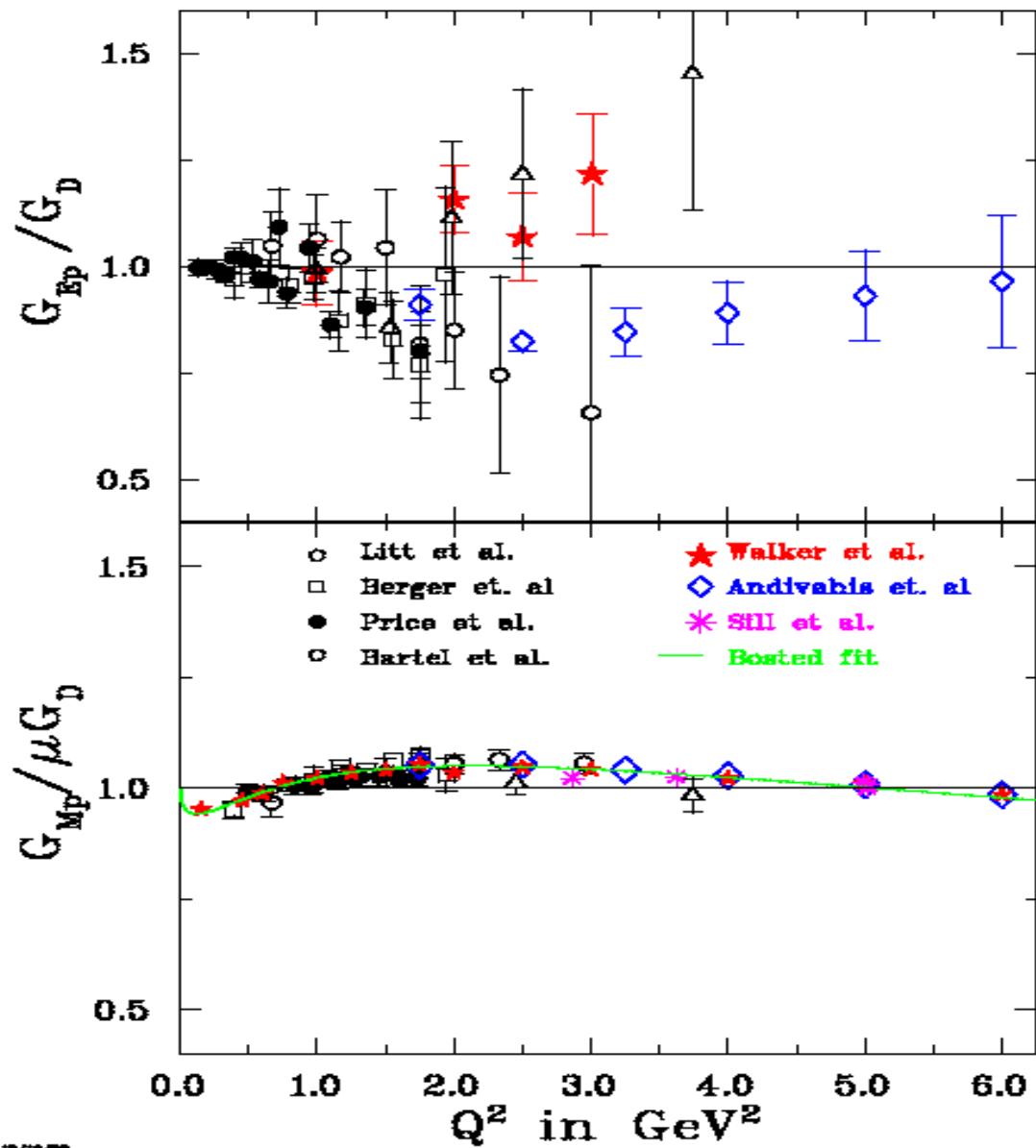
$$\sigma_R = \frac{d\sigma}{d\Omega} \frac{(1 + \tau)\varepsilon}{\sigma_{ns} \tau} = \frac{\varepsilon}{\tau} \mathbf{G}_{Ep}^2(Q^2) + \mathbf{G}_{Mp}^2(Q^2)$$

$$\varepsilon = \{1 + 2(1 + \tau) \tan^2(\vartheta_e / 2)\}^{-1}$$

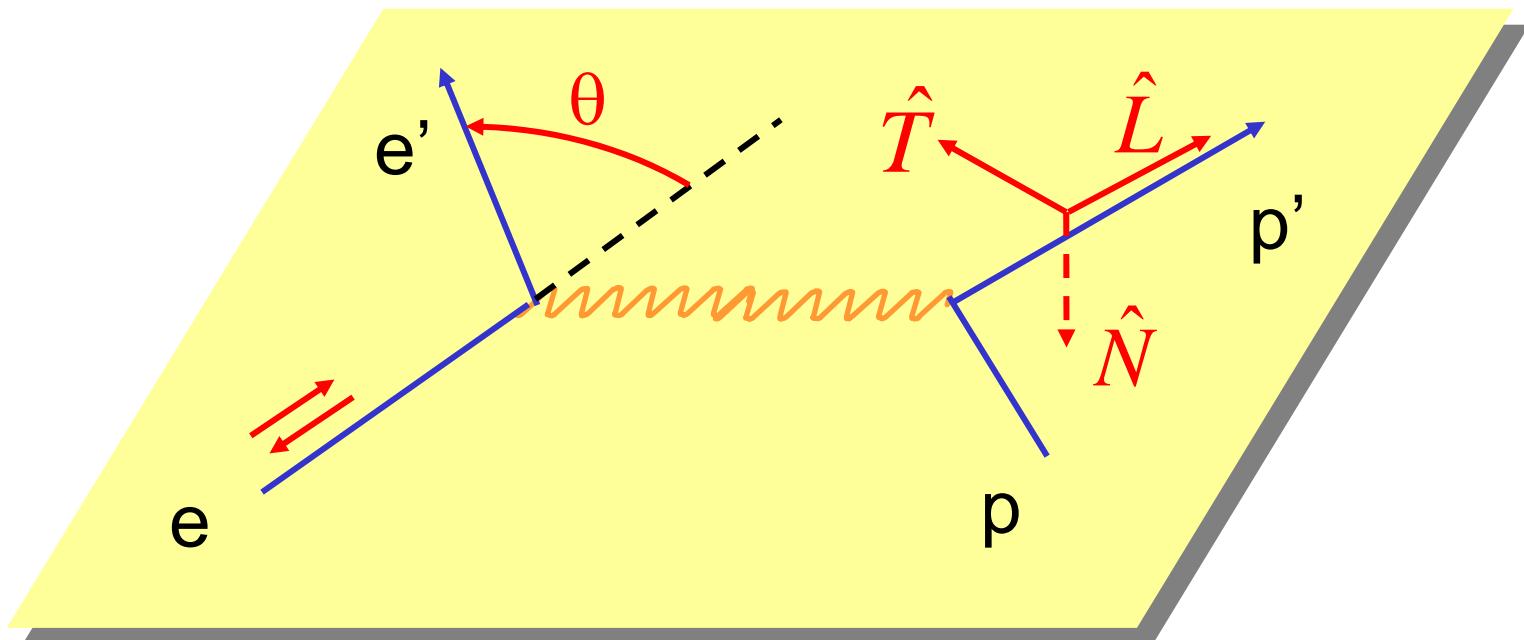
Proton Form Factors pre-1998

$$G_{Ep} \sim \frac{G_{Mp}}{\mu} \sim G_D$$

$$G_D = \frac{1}{\left(1 + \frac{Q^2}{0.71}\right)^2}$$

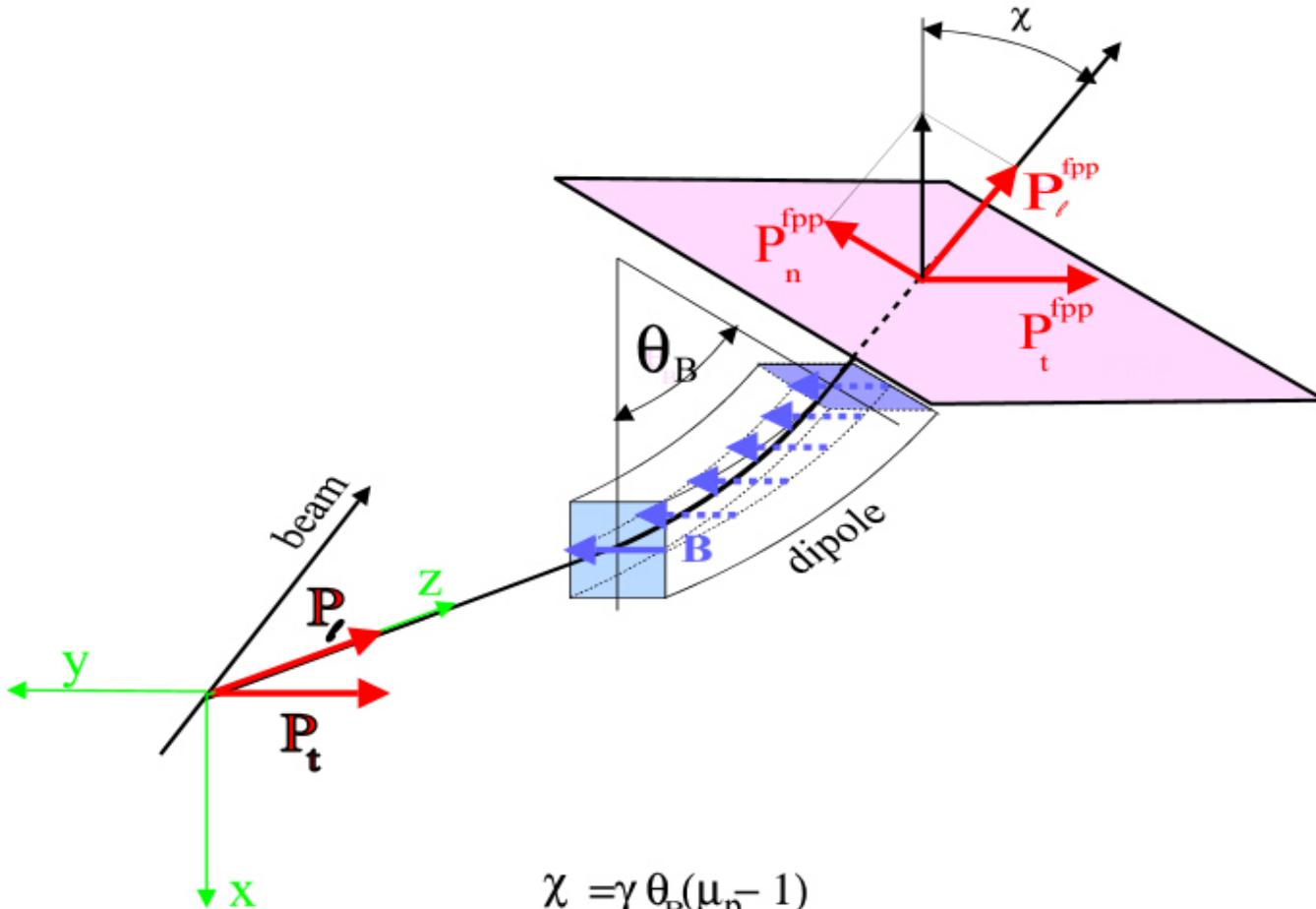


Spin transfer reaction $\vec{e} \vec{p} \rightarrow e \vec{p}'$



$$\frac{G_{Ep}}{G_{Mp}} = - \frac{P_t}{P_l} \frac{E_e + E_{e'}}{2M} \tan(\vartheta_e/2)$$

Transport through magnet

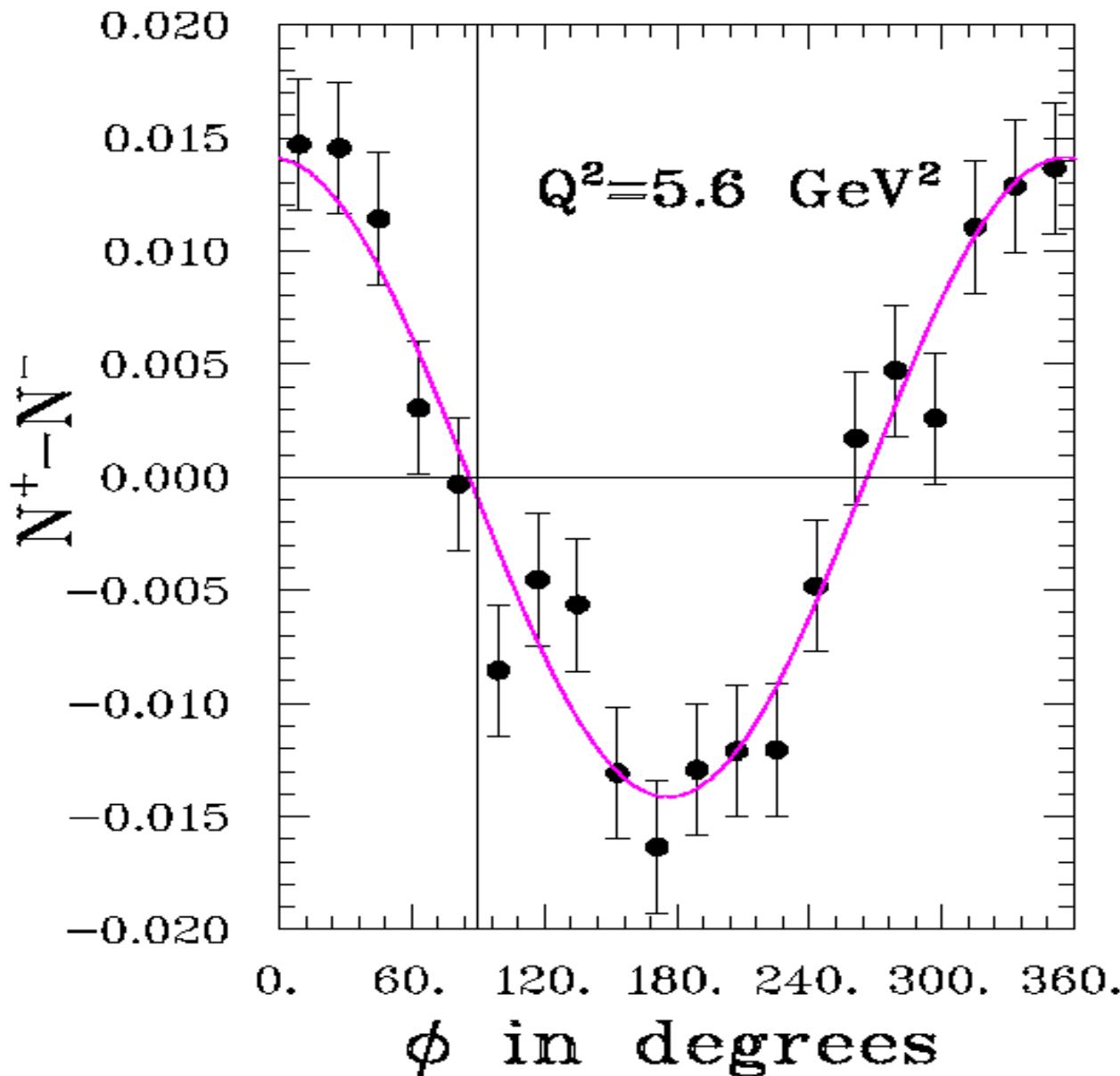


$$\chi = \gamma \theta_B (\mu_p - 1)$$

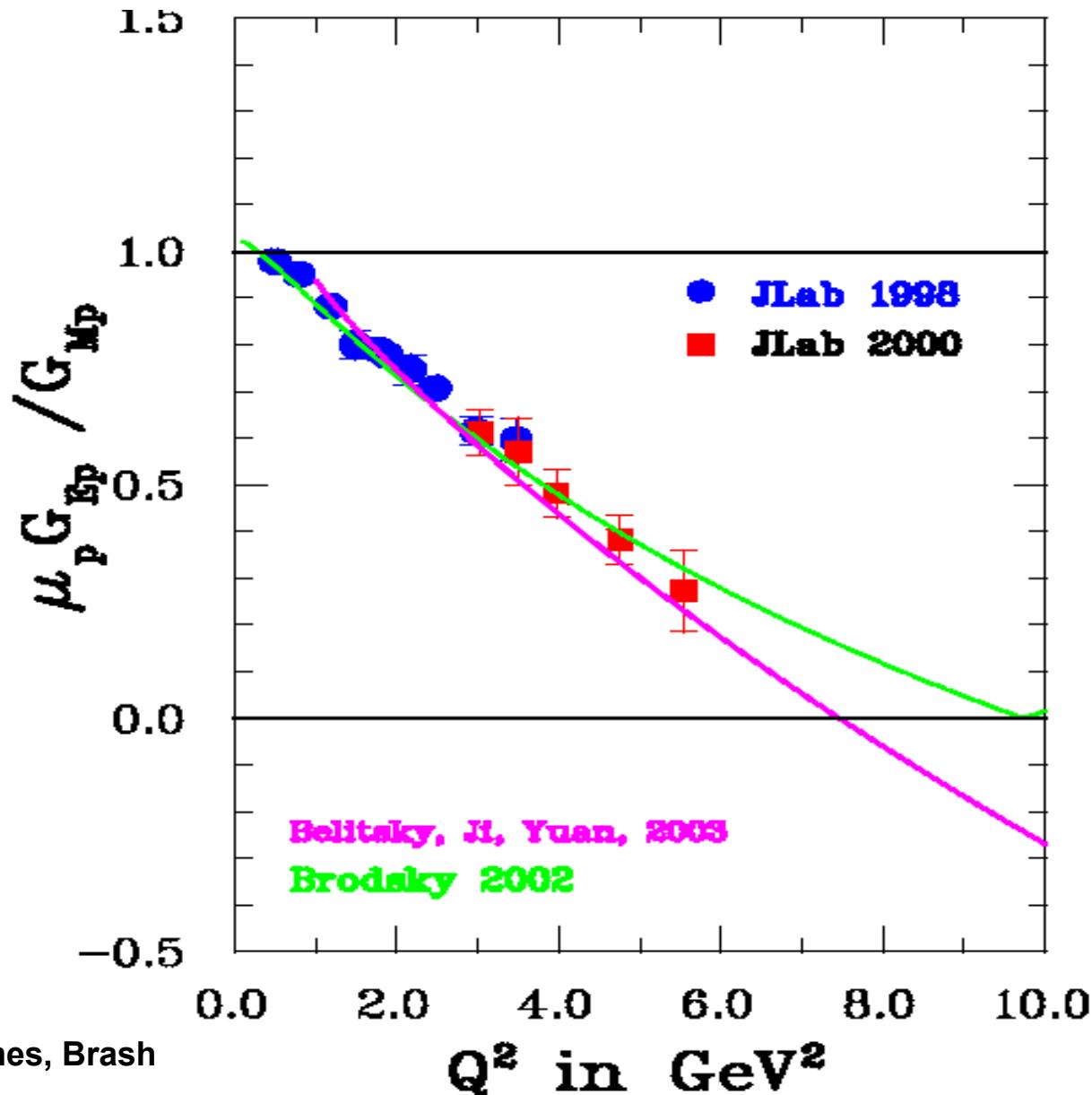
$$P_t^{\text{fpp}} = P_t$$

$$P_\ell^{\text{fpp}} = P_\ell \sin \chi$$

Azimuthal asymmetry in the polarimeter



G_{Ep} from polarization transfer

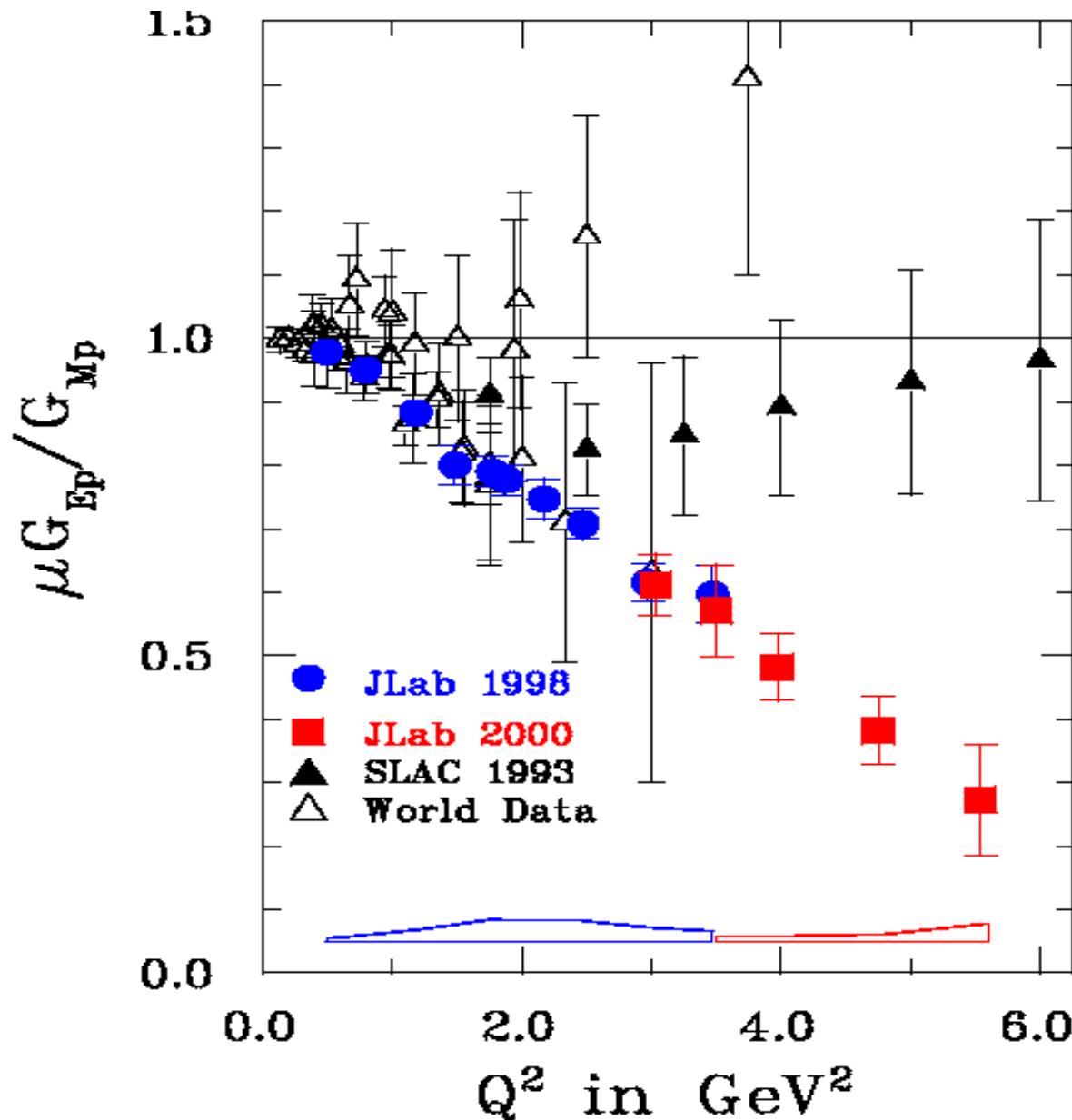


E93-027, E99-007

Perdrisat, Punjabi, Jones, Brash



World data for G_{Ep}



Interpretation of new data

$F_2(Q^2)$ is a spin-flip transition

$$\langle P'_\uparrow | J^\mu | P_\downarrow \rangle \sim F_2(Q^2) \bar{u}_\uparrow(P') \frac{i\sigma^{\mu\nu} q_\alpha}{2M} u_\downarrow(P)$$

In the absence of quark angular momentum

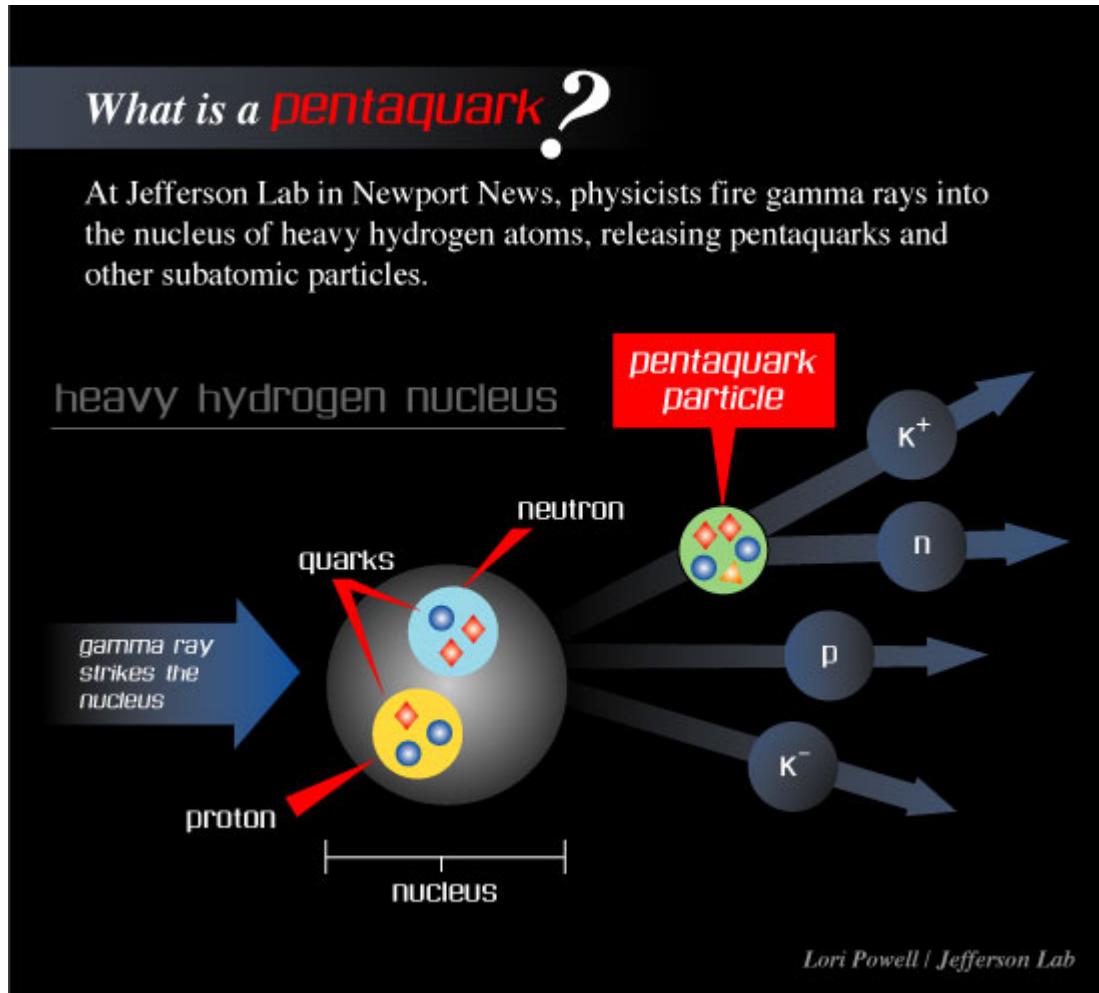
$$Q^2 \frac{F_2}{F_1} \sim m_q M \xrightarrow{m_q \rightarrow 0} 0$$

Quark orbital angular momentum
essential to describe data

$$Q \frac{F_2}{F_1} \sim \frac{1}{Q} \log^2 \left(\frac{Q^2}{\Lambda^2} \right) \sim const$$

Pentaquark: Baryon with five quarks

Goal: Determine quark content of colorless hadrons

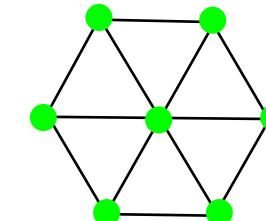


Expectation from the quark model is that the properties of baryons are determined by three valence quarks (qqq)

Hadron multiplets

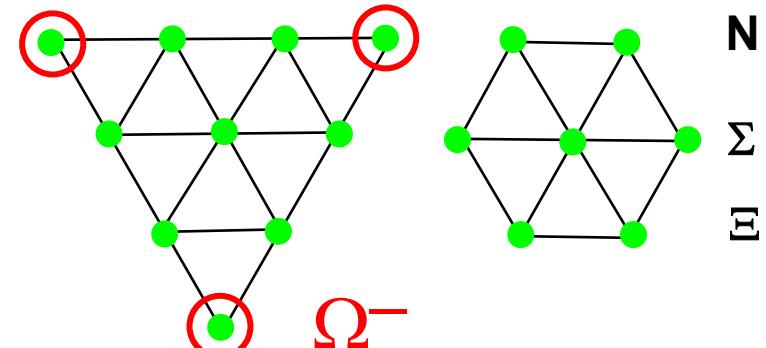
Mesons $q\bar{q}$

$$3 \otimes \bar{3} = 8 \oplus 1$$



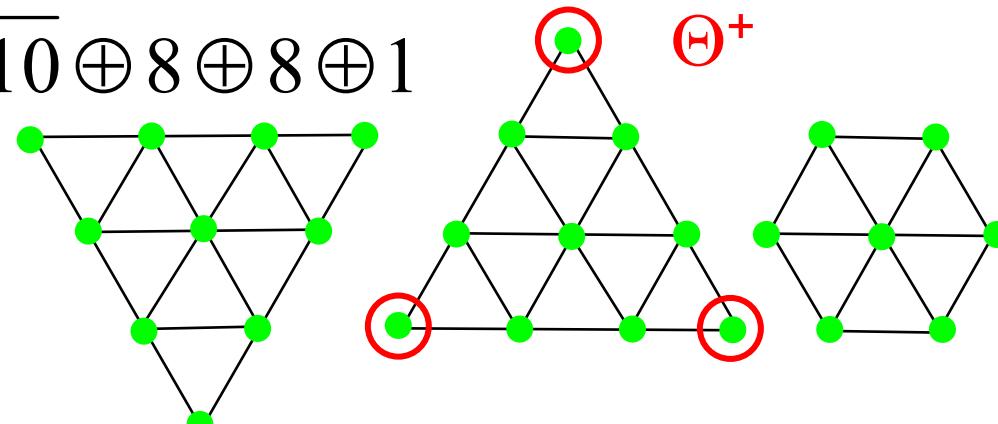
Baryons qqq

$$3 \otimes 3 \otimes 3 = 10 \oplus 8 \oplus 8 \oplus 1$$



Baryons built from meson-baryon basis

$$8 \otimes 8 = 27 \oplus 10 \oplus \bar{10} \oplus 8 \oplus 8 \oplus 1$$



Production and decay of $\Omega^- \rightarrow \Xi^0 \pi^-$

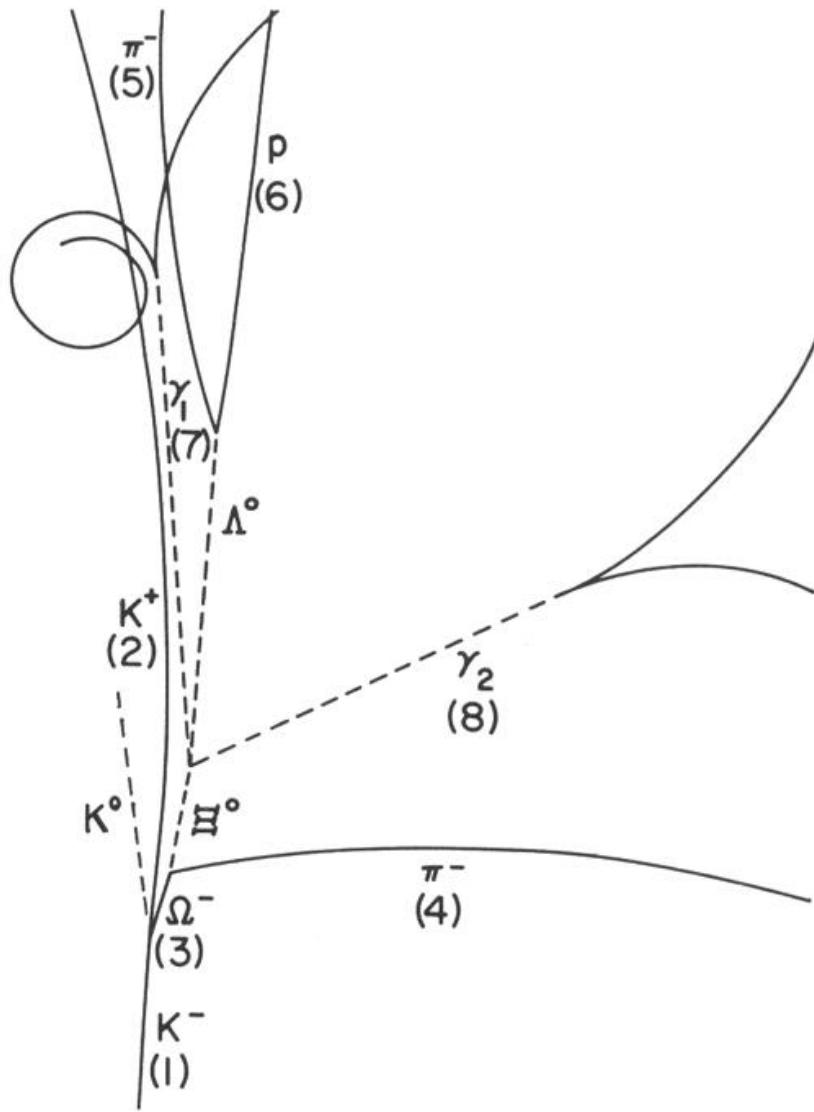
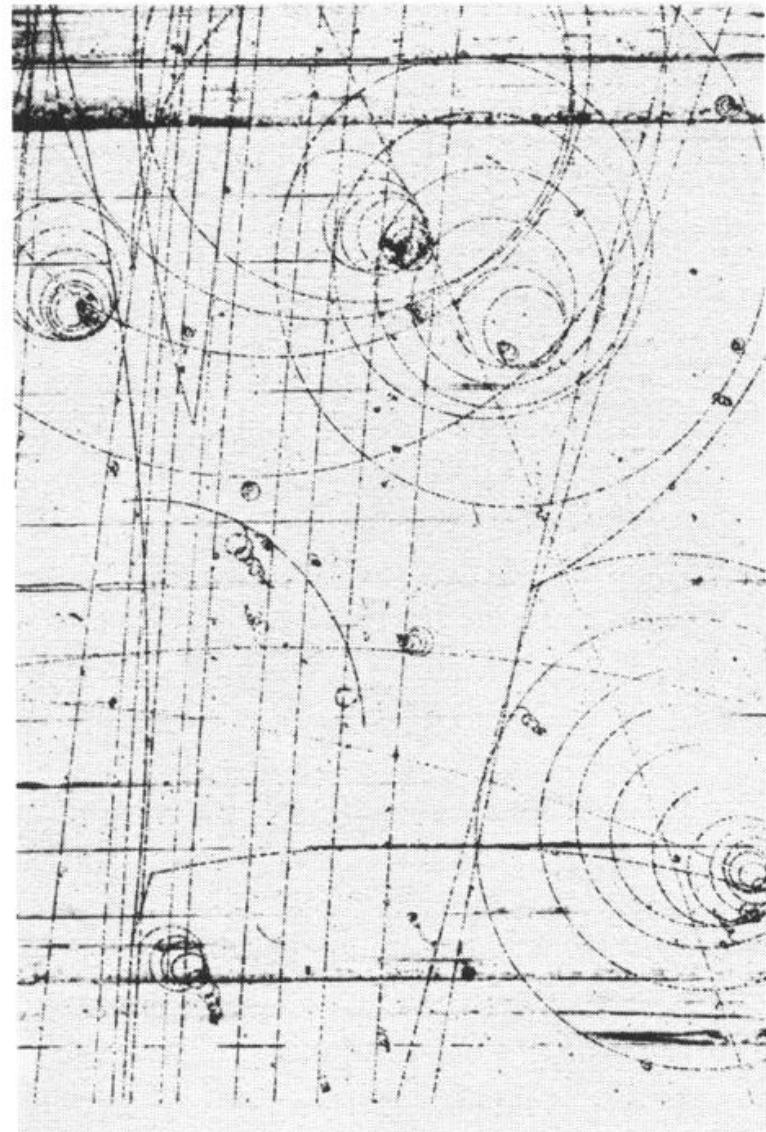
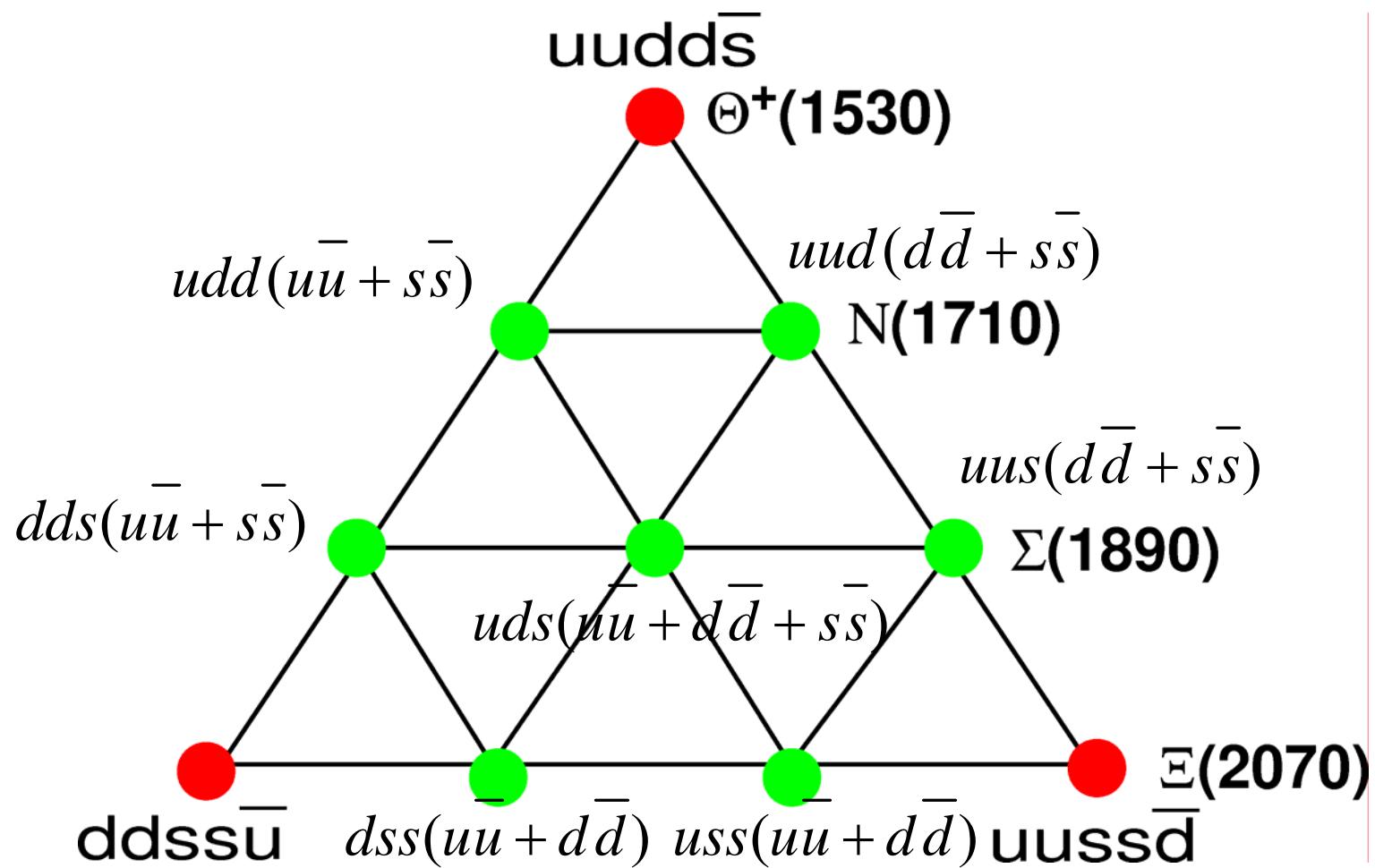


FIG. 2. Photograph and line diagram of event showing decay of Ω^- .

What are pentaquarks?

- Minimum quark content is 5-quarks.
 - Anti-quark has different flavor than any of 4-quarks ($qqqq\bar{Q}$).
 - Quantum numbers can not be defined by 3-quarks.
-
- General idea of a five-quark states has been around since late 60's.
 - However, searches did not give any conclusive results.
 - PDG dropped the discussion on pentaquark searches after 1988.

The Anti-decuplet predicted by Diakonov et al.



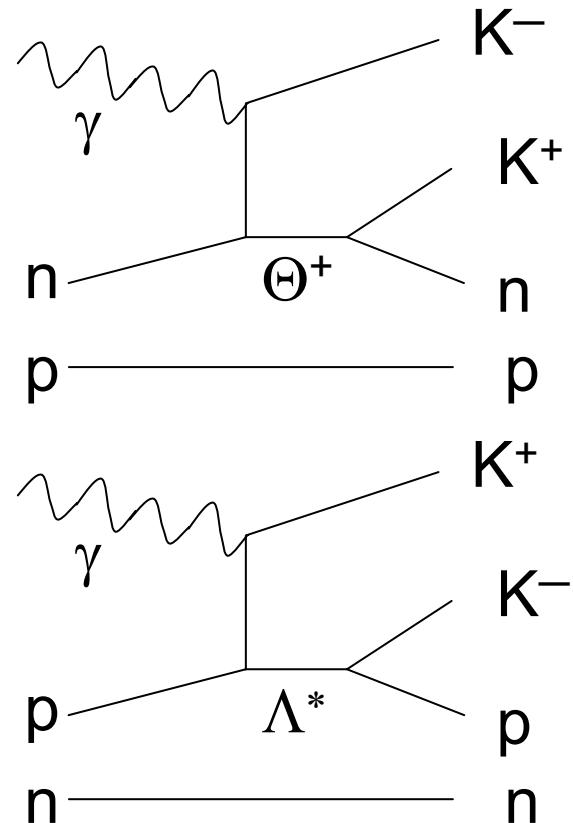
Reactions on deuterium

$$\gamma n(p) \rightarrow \Theta^+ K^-(p)$$

$$\Theta^+ \rightarrow K^+ n$$

$$\gamma p(n) \rightarrow \Lambda^*(1520) K^+(n)$$

$$\Lambda^*(1520) \rightarrow K^- p$$



$$\gamma N \rightarrow \phi(1020) N \rightarrow K^+ K^- N$$

CEBAF Large Acceptance Spectrometer

Torus magnet

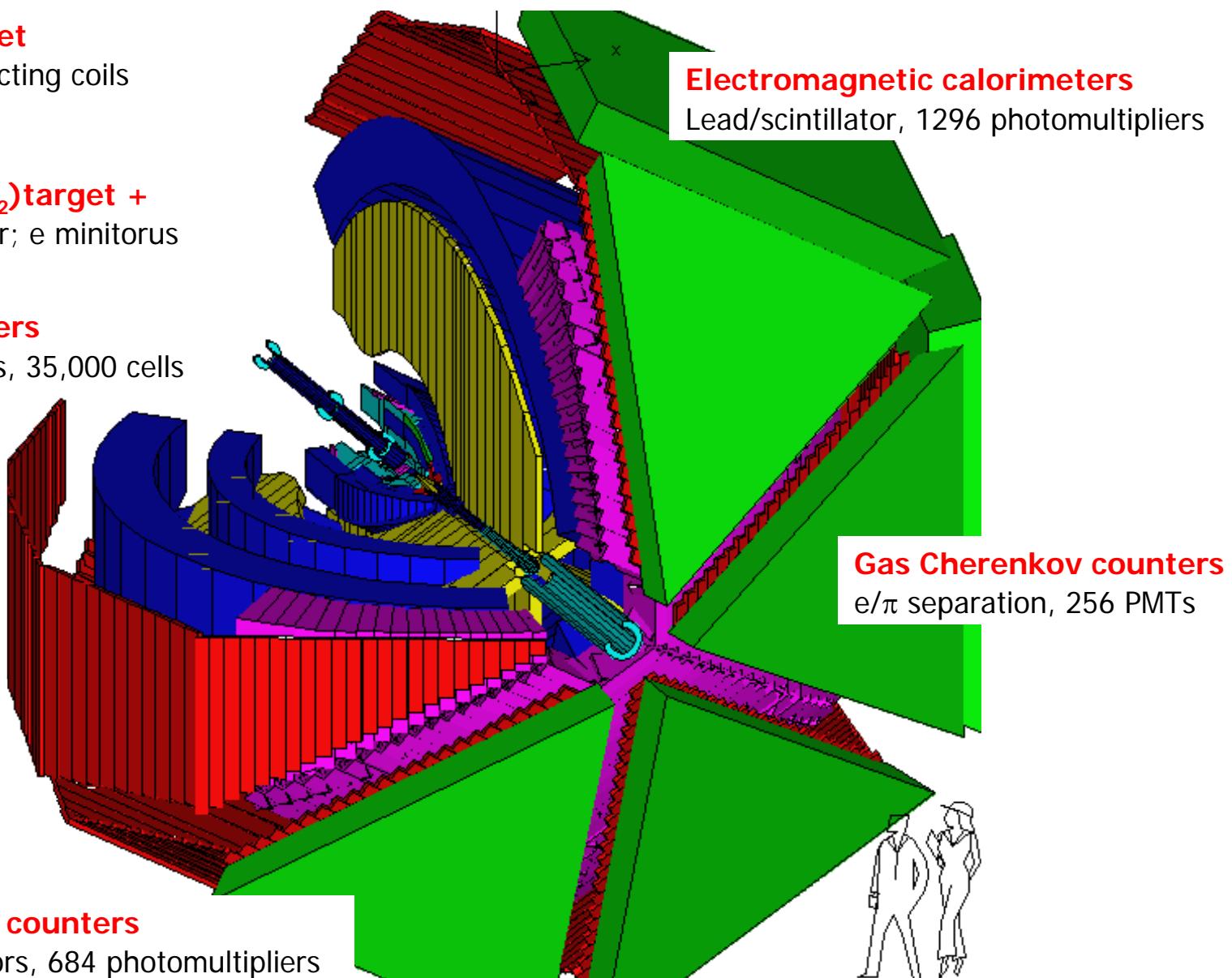
6 superconducting coils

Liquid D₂ (H₂) target +

γ start counter; e minitorus

Drift chambers

argon/CO₂ gas, 35,000 cells



Time-of-flight counters

plastic scintillators, 684 photomultipliers

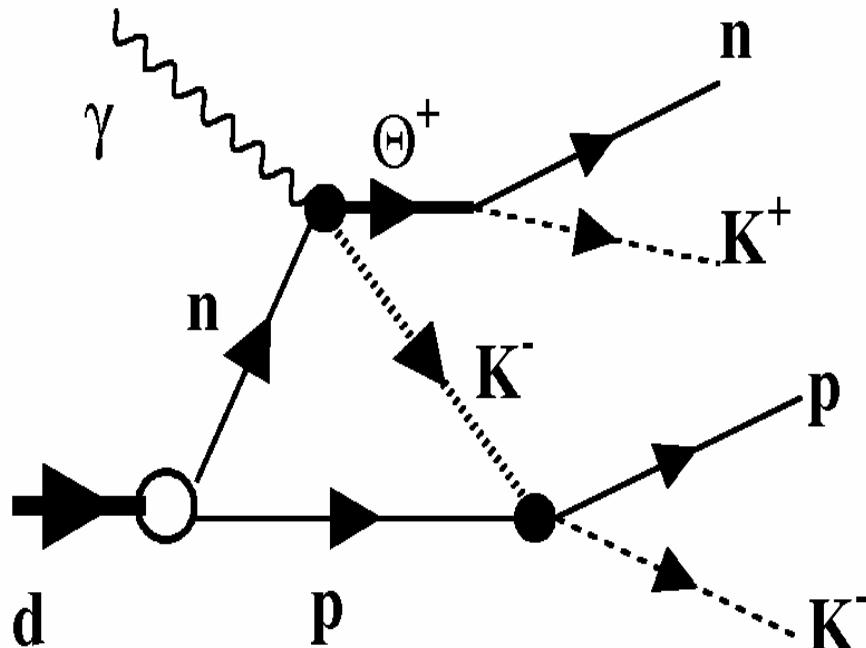
Electromagnetic calorimeters

Lead/scintillator, 1296 photomultipliers

Gas Cherenkov counters

e/ π separation, 256 PMTs

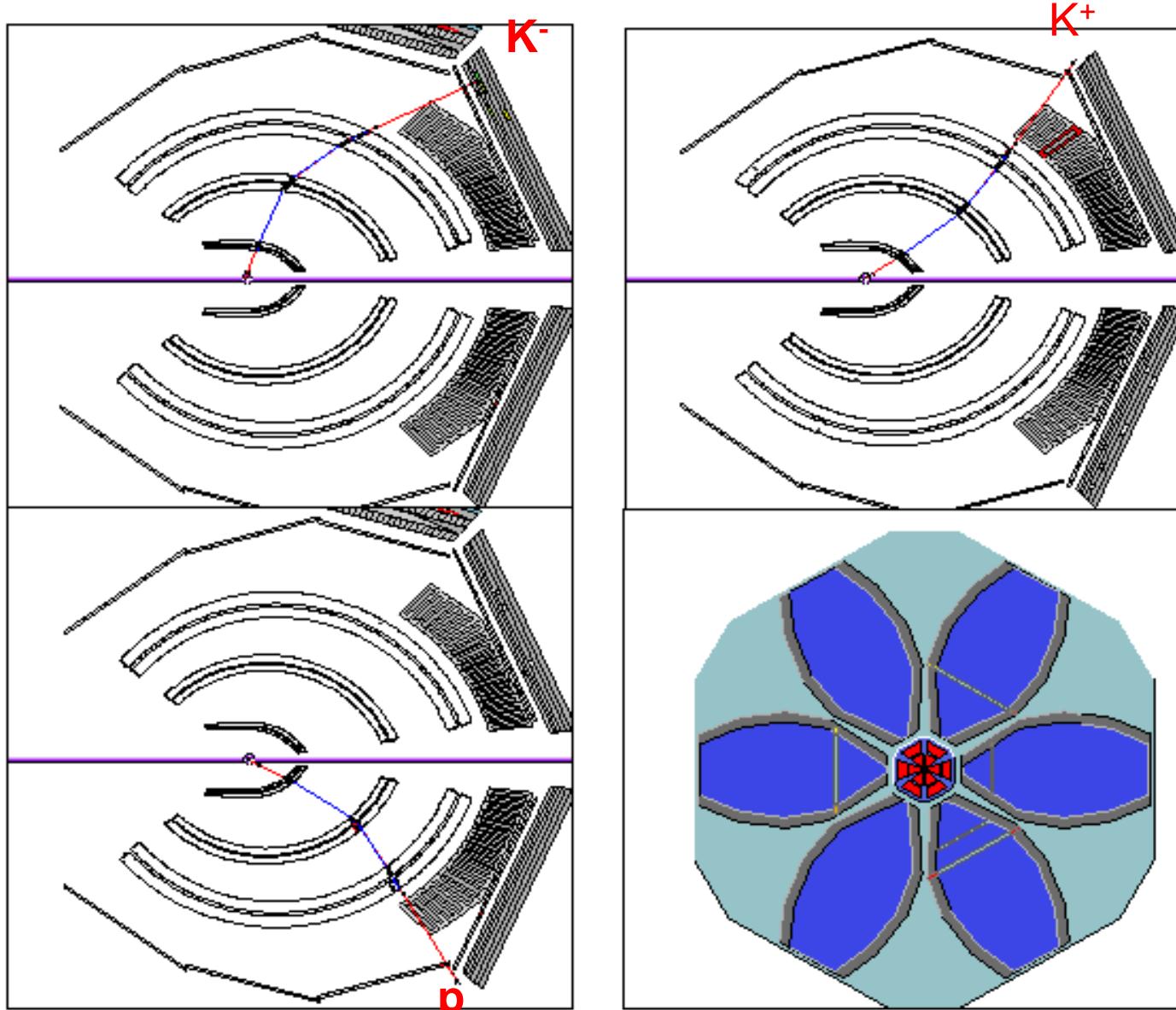
Exclusive measurement using γd reactions



CLAS Collaboration (S. Stepanyan, K. Hicks, *et al.*), hep-ex/0307018

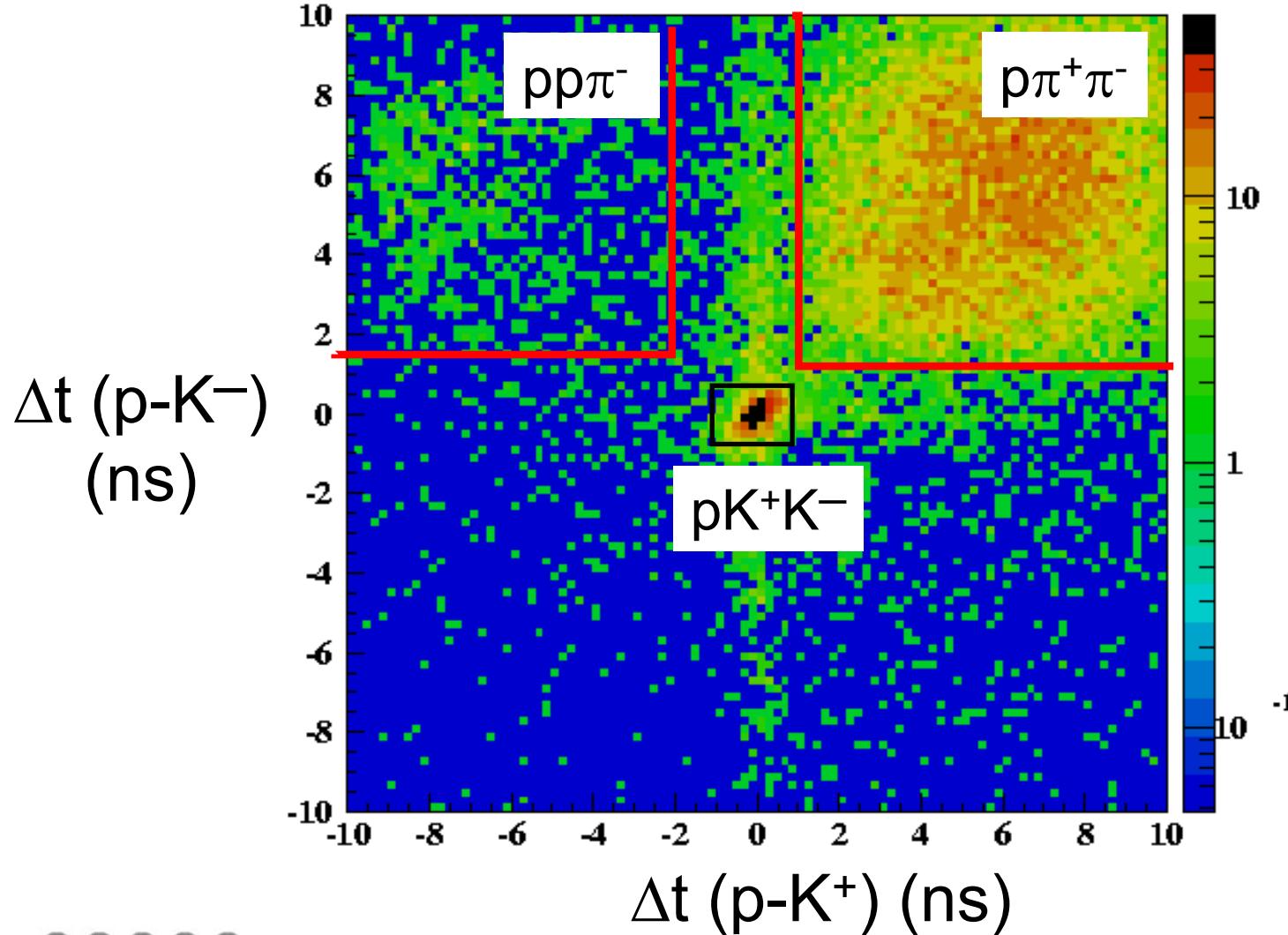
- Requires FSI - both nucleons involved
 - No Fermi motion correction necessary
 - FSI puts K^- at larger lab angles: better CLAS acceptance
 - FSI not rare: in ~50% of $\Lambda^*(1520)$ events both nucleons detected with $p > 0.15 \text{ GeV}/c$

$\gamma d \rightarrow p K^+K^- (n)$ in CLAS



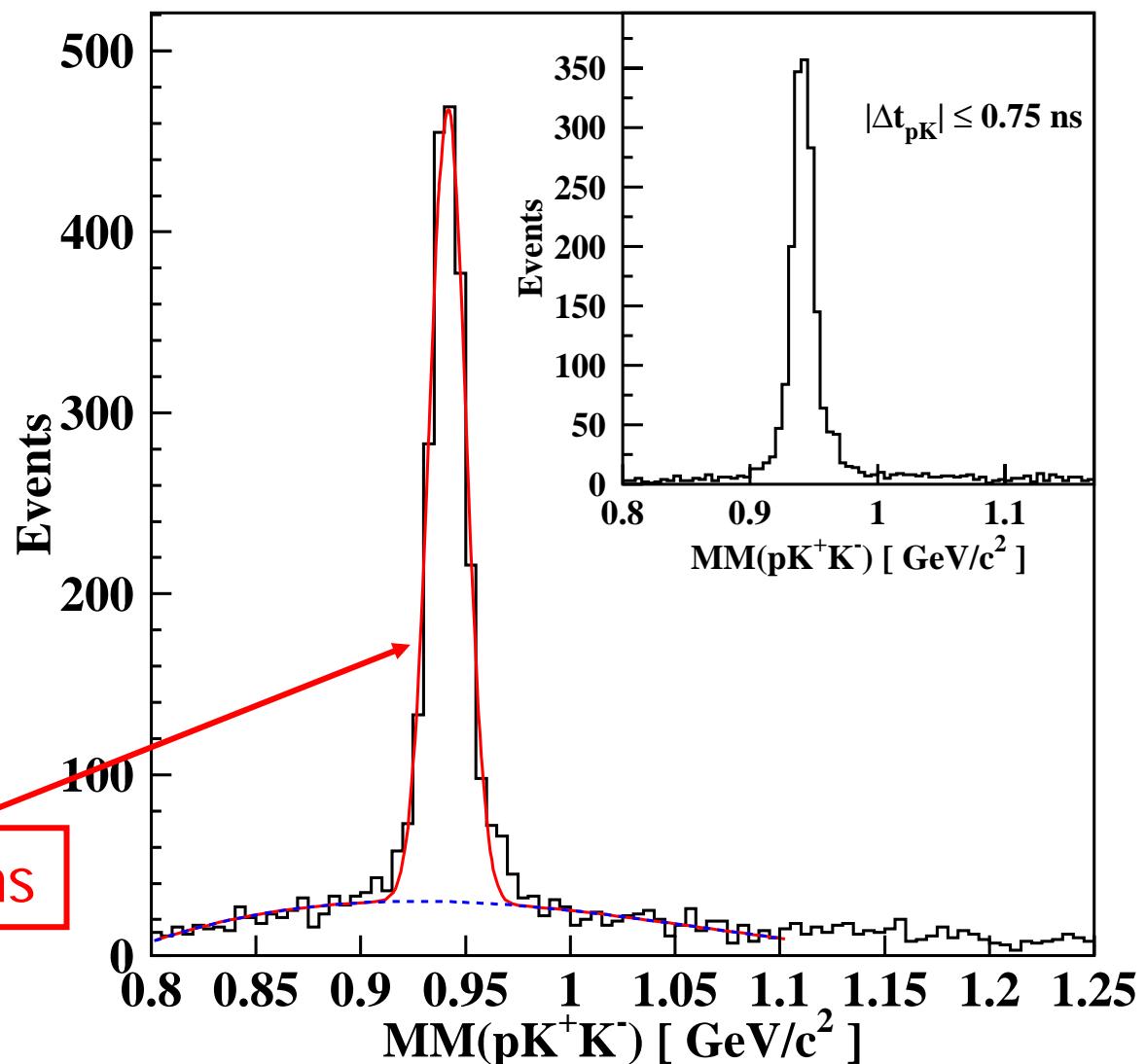
Kaon times relative to proton

$$\Delta t_K = t - \frac{R}{\beta_c \cdot c}; \beta_c = \frac{p}{\sqrt{p^2 + m_K^2}}$$

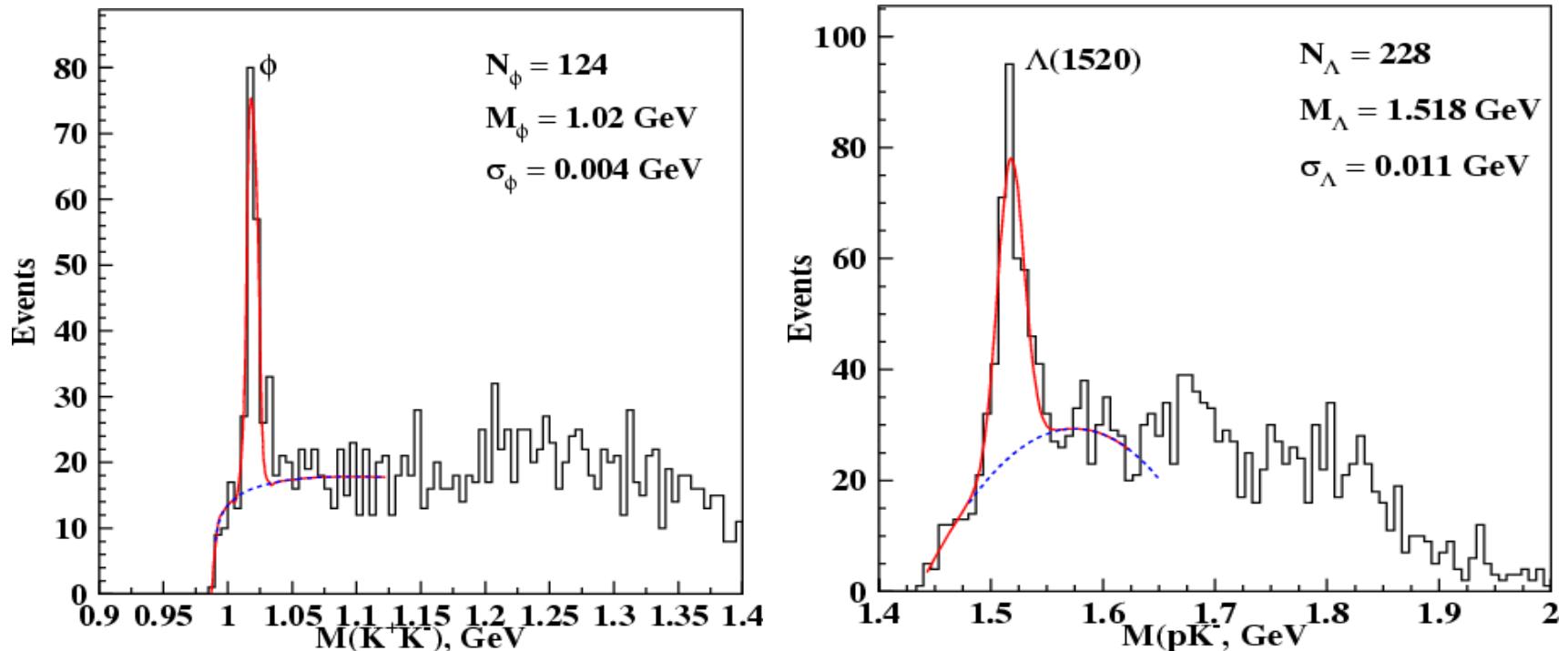


Reaction $\gamma d \rightarrow p K^+ K^- (n)$

- Clear peak at neutron mass.
- 15% non pKK events within $\pm 3\sigma$ of the peak.
- Almost no background under the neutron peak after event selection with tight timing cut.

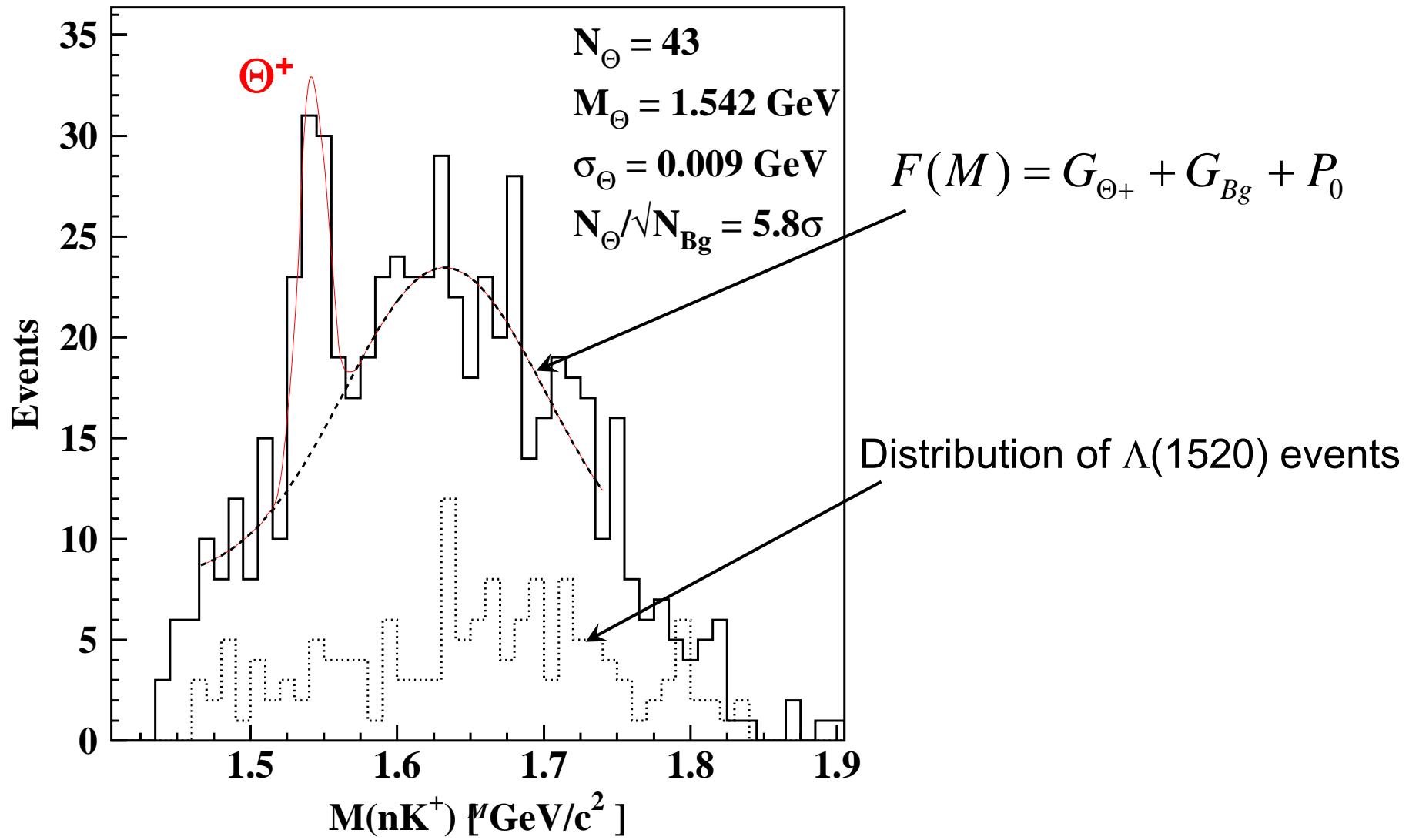


Identification of known resonances



- Remove events with $IM(K^+K^-) \rightarrow \phi(1020)$ by $IM > 1.07 \text{ GeV}$
- Remove events with $IM(pK^-) \rightarrow \Lambda(1520)$
- Limit K^+ momentum due to $\gamma d \rightarrow p K^- \Theta^+$ phase space $p_{K^+} < 1.0 \text{ GeV}/c$
- C. Meyer (CLAS note 03-009):** checked narrow structure impossible in $\gamma d \Rightarrow K^+ Y^* N \Rightarrow K^+ (K^- N) N, + KN$ rescattering

nK^+ invariant mass distribution



Θ^+ : experimental status

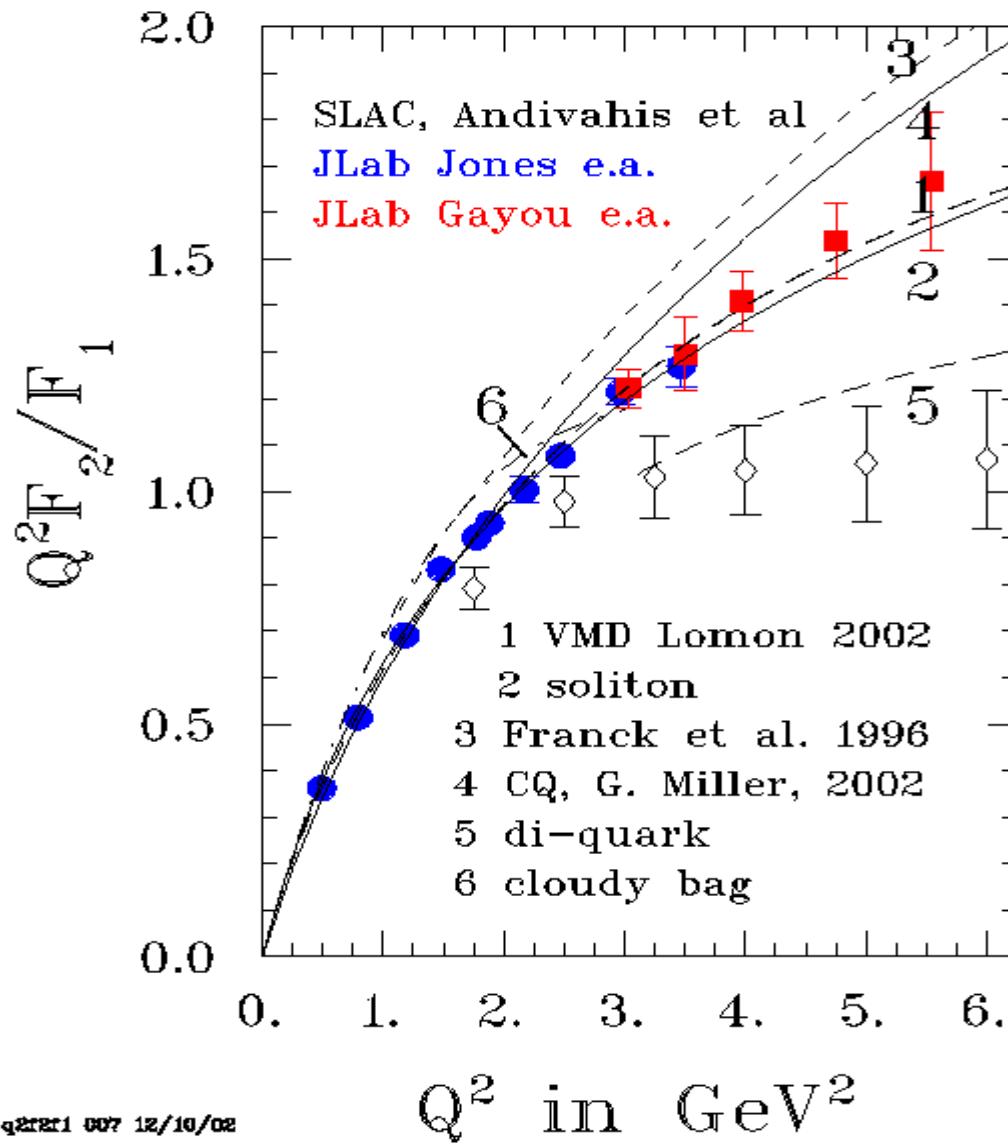
- Experimental evidence for Θ^+ have been reported at four laboratories.
 - LEPS collaboration at Spring-8 (Japan), January 2003 - peak in the invariant mass of the nK^+ at 1.54 GeV with statistical significance of 4.6σ .
 - DIANA collaboration at ITEP (Moscow), April 2003 – peak in the invariant mass of pK^0 at 1.538 GeV, statistical significance 4.4σ .
 - CLAS collaboration at JLAB, July 2003 – peak in the invariant mass of the nK^+ at 1.542 GeV, statistical significance 5.3σ .
 - SAPHIR collaboration at ELSA (Bonn), August 2003 – peak in the invariant mass of the nK^+ at 1.54 GeV, statistical significance 4.8σ .
- All experiments observe a narrow width.
- Spin, isospin and parity not yet established.
- Subject of intense interest and research.
 - Penta-Quark 2003 Workshop at JLab in November.

Summary

- We have presented two examples which highlight the physics program at Jefferson Lab.
- The electromagnetic interaction can be used to probe deep into the structure of nucleons.
 - From measurements of G_{Ep} up to a $Q^2 = 5.6 \text{ GeV}^2$ we have gained new insights into the shape of the proton.
 - Orbital angular momentum of quarks is a key ingredient in our understanding of proton structure.
- A key question in non-perturbative QCD is the structure of hadrons
 - We have presented evidence for an exotic baryon with $S = +1$, which would have a minimal quark content of five quarks ($uudd\bar{s}$).
 - This baryon represents a new class of colorless hadrons.



Scaled F_2/F_1 ratio



q20201 007 12/10/02

Q^2 in GeV^2