



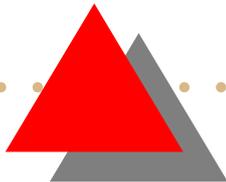
*Results from G0*

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for the G0 Collaboration

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Caltech, CMU, William & Mary, Hampton, IPN-Orsay, ISN-Grenoble, JLab, Kentucky,  
LaTech, NMSU, TRIUMF, U Conn, UIUC, U Manitoba, U Maryland, U Mass UNBC,  
VPI, Yerevan



# Physics Motivation



$s$  quark distribution is interesting since it exclusively comes from  $s\bar{s}$  sea

- The proton electromagnetic form factors is weighted sum of the quark form factors

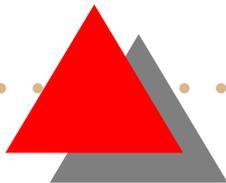
$$G_{E(M)}^{\gamma,p} = \frac{2}{3}G_{E(M)}^{u,p} - \frac{1}{3}G_{E(M)}^{d,p} - \frac{1}{3}G_{E(M)}^{s,p}$$

- Measuring  $G_{E(M)}^{\gamma,p}$  not enough to separate the quark form factors

$$G_{E(M)}^{Z,p} = (1 - \frac{8}{3} \sin^2 \theta_W)G_{E(M)}^{u,p} + (-1 + \frac{4}{3} \sin^2 \theta_W)G_{E(M)}^{d,p} + (-1 + \frac{4}{3} \sin^2 \theta_W)G_{E(M)}^{s,p}$$

- Measurement of  $G_{E(M)}^{Z,p}$  add 2 new eqs. Still not enough, but one can extract

$$\begin{aligned} G_{E(M)}^0 &= \frac{1}{3}(G_{E(M)}^{u,p} + G_{E(M)}^{d,p} + G_{E(M)}^{s,p}) \\ &= (2 - 4 \sin^2 \theta_W)G_{E(M)}^{\gamma,p} - G_{E(M)}^{Z,p} \end{aligned}$$



# Physics Motivation

- assuming that proton and neutron differ only by exchanging  $u$  and  $d$

$$G_{E(M)}^{u,n} = G_{E(M)}^{d,p} \quad G_{E(M)}^{d,n} = G_{E(M)}^{u,p} \quad G_{E(M)}^{s,n} = G_{E(M)}^{s,p}$$

$$G_{E(M)}^{\gamma,n} = \frac{2}{3}G_{E(M)}^{d,p} - \frac{1}{3}G_{E(M)}^{u,p} - \frac{1}{3}G_{E(M)}^{s,p}$$

- quark form factors can be isolated

$$G_{E(M)}^{s,p} = (1 - 4 \sin^2 \theta_W)G_{E(M)}^{\gamma,p} - G_{E(M)}^{\gamma,n} - G_{E(M)}^{Z,p}$$

From measuring  $G_{E(M)}^{Z,p}$  and using the well known  $G_{E(M)}^{\gamma,p}$  and  $G_{E(M)}^{\gamma,n}$  one can extract individual quark form factors



# How to measure $G_{E(M)}^{Z,p}$ ?

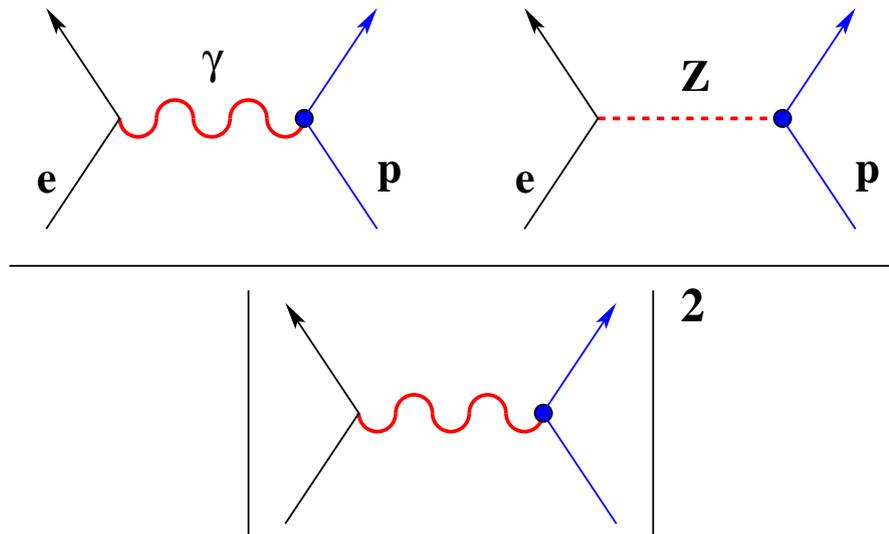
PV asymmetry for elastic electron-proton scattering  $A = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L}$

$$A = \frac{G_F Q^2}{4\sqrt{2}\pi\alpha} \frac{\epsilon G_E^{\gamma,p} G_E^{Z,p} + \tau G_M^{\gamma,p} G_M^{Z,p} - (1 - 4\sin^2\theta_W)\epsilon' G_M^{\gamma,p} G_A^e}{\epsilon (G_E^{\gamma,p})^2 + \tau (G_M^{\gamma,p})^2}$$

$$\tau = \frac{Q^2}{4M_N^2}$$

$$\epsilon = (1 + 2(1 + \tau) \tan^2(\theta_2/2))^{-1}$$

$$\epsilon' = \sqrt{\tau(1 + \tau)(1 - \epsilon^2)}$$



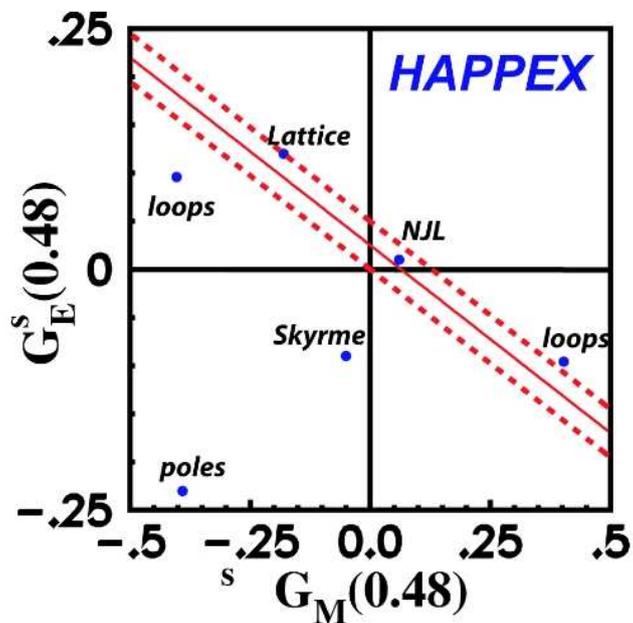
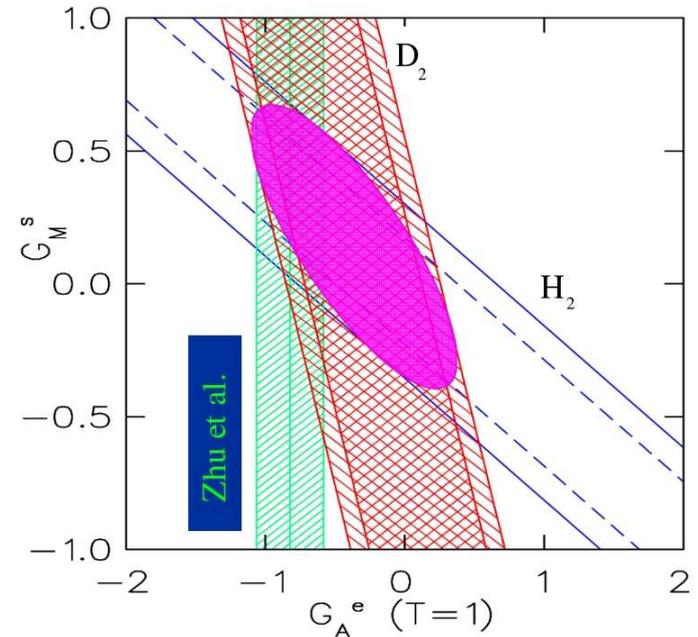
# Previous measurements

## SAMPLE at MIT-Bates

$$\vec{e} + p \quad A_p = -4.92 \pm 0.61 \pm 0.73 \text{ ppm}$$

$$\vec{e} + d \quad A_d = -7.55 \pm 0.70 \pm 0.60 \text{ ppm}$$

$$G_M^s(Q^2 = 0.1 \text{ GeV}^2) = 0.14 \pm 0.35 \pm 0.40$$



## HAPPEX at JLAB

$$\vec{e} + p \quad A_p = -15.05 \pm 0.98 \pm 0.56 \text{ ppm}$$

$$G_E^s + 0.39G_M^s = 0.025 \pm 0.020 \pm 0.014$$

$$\text{at } Q^2 = 0.48 \text{ GeV}^2$$

# G0 experiment



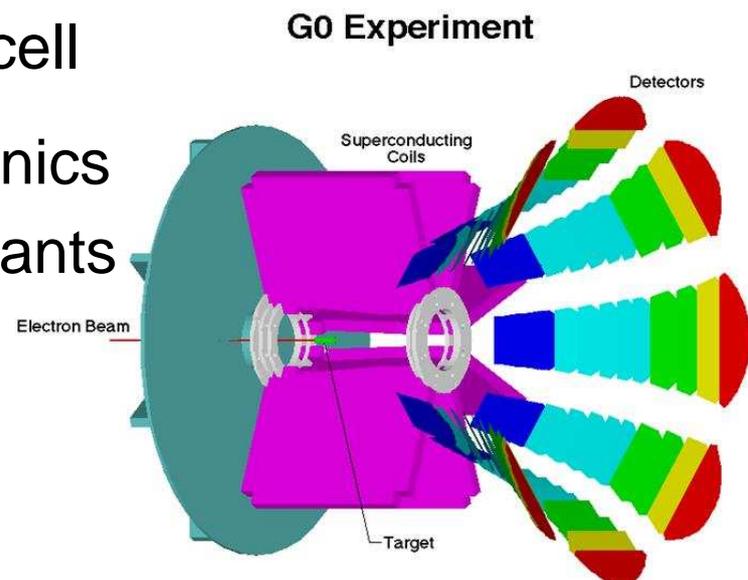
## Goals of the G0 experiment:

- measure the  $A \sim -3$  to  $-40$  ppm with precision of  $\sim 5\%$ 
  - need  $10^{13}$  to  $10^{14}$  events
  - high polarization, high current
  - high count rate detectors/electronics
  - small helicity correlated beam properties
  - control over sources of false asymmetries (deadtime, background, ..)
- Separate  $G_E^s$  and  $G_M^s$ 
  - Forward angle elastic  $\vec{e} + p$  - Full range of  $Q^2=0.1-1.0$  GeV<sup>2</sup>/c<sup>2</sup>
  - Backward angle elastic  $\vec{e} + p$  -  $Q^2=0.3, 0.5, 0.8$  GeV<sup>2</sup>/c<sup>2</sup>
  - Backward angle quasi-elastic  $\vec{e} + d$  -  $Q^2=0.3, 0.5, 0.8$  GeV<sup>2</sup>/c<sup>2</sup>

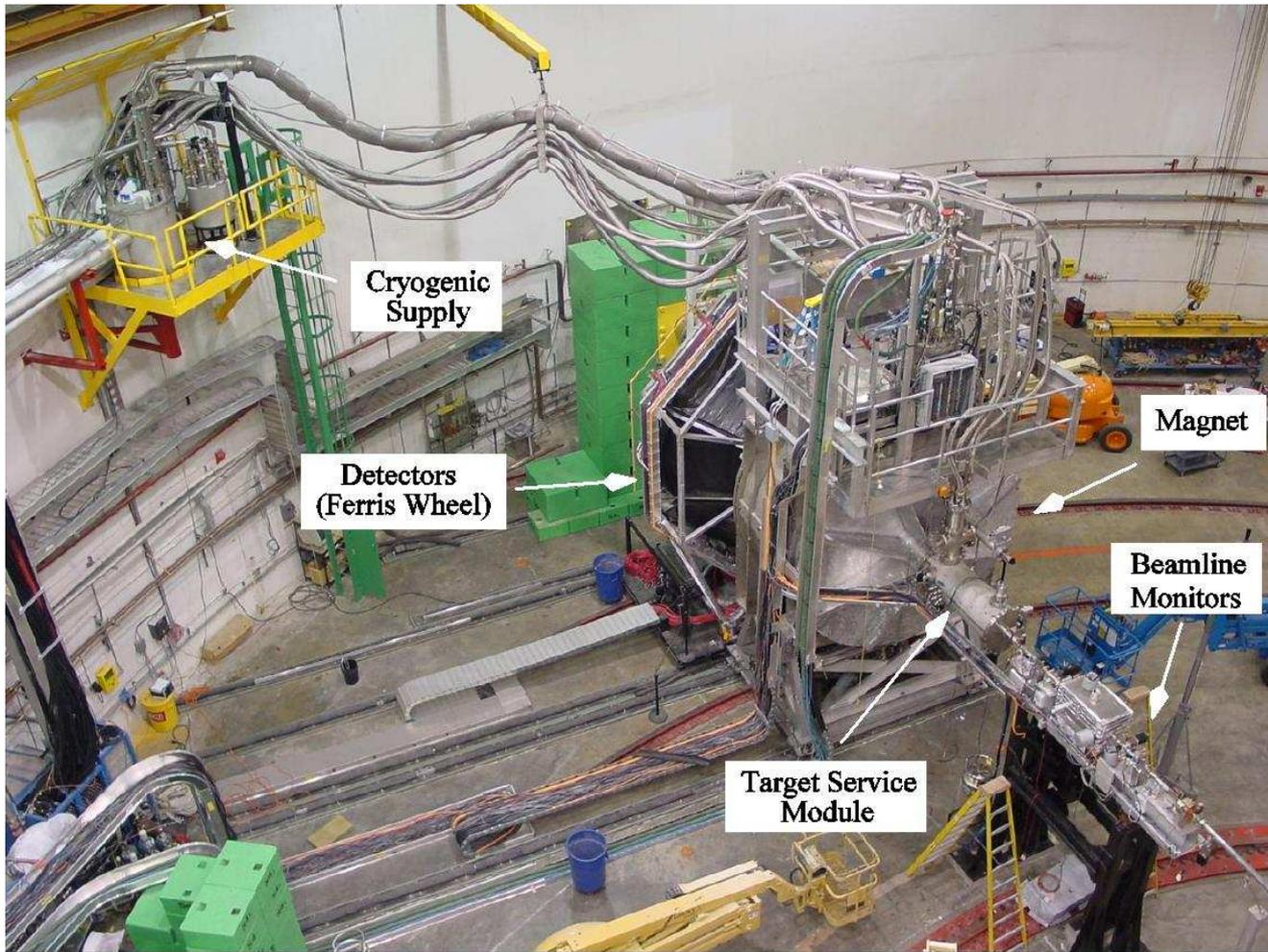


# Forward angle set-up overview

- superconducting toroidal magnet
- segmented large area scintillation detector array
- custom and commercial high rate electronics
- 3 GeV polarized beam in 32  $ns$  pulses
- High power 20 cm LH<sub>2</sub> target cell
- Different hardware and electronics (NA/French) for alternating octants



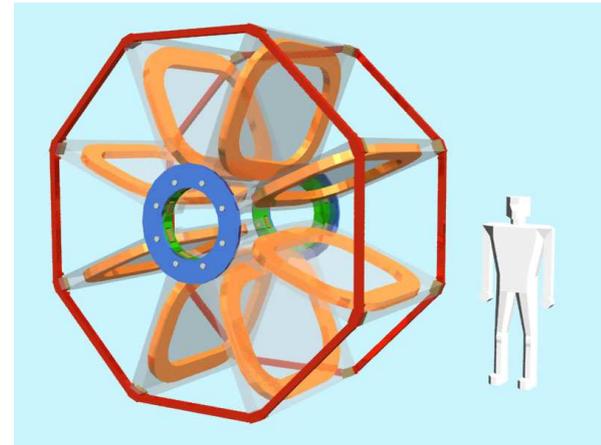
# *G0 experimental set-up*



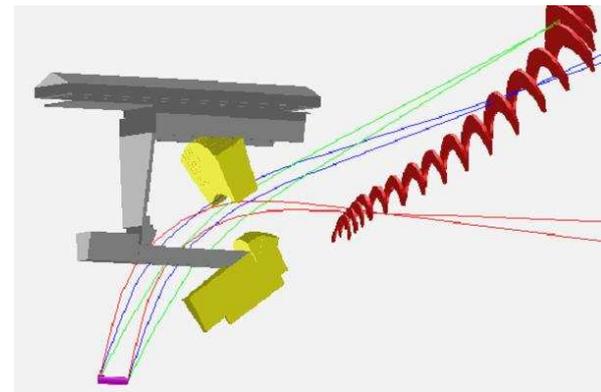
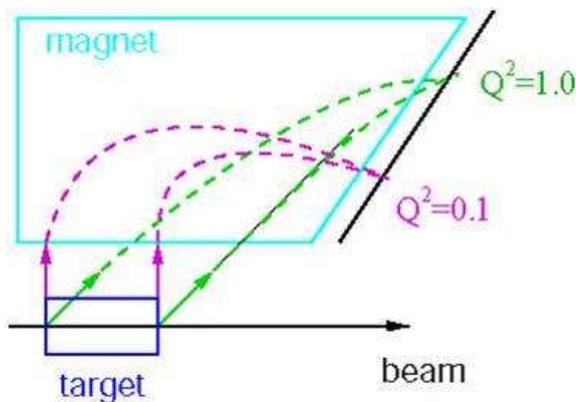
# G0 magnet

## Superconducting toroidal magnet

- 8 coils
- $\int Bdl = 1.6 \text{ Tm}$
- $35^\circ < \theta_{bend} < 86^\circ$
- $\phi$  acceptance  $\sim 0.44$
- Ran at full design current (5000 A)



Magnet "sorts" proton by  $Q^2$  in focal plane detector

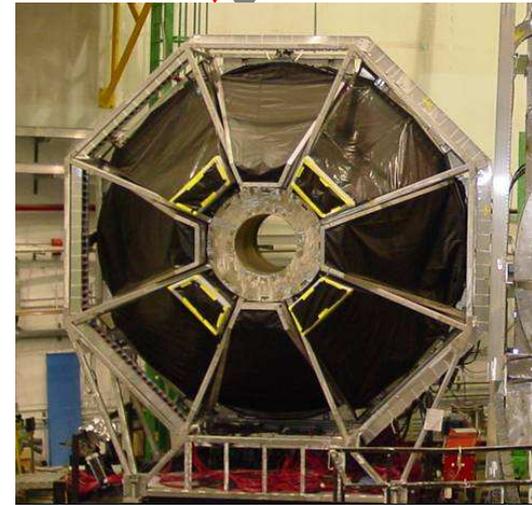


# Focal plane detectors

16 pairs of arc-shaped scintillators

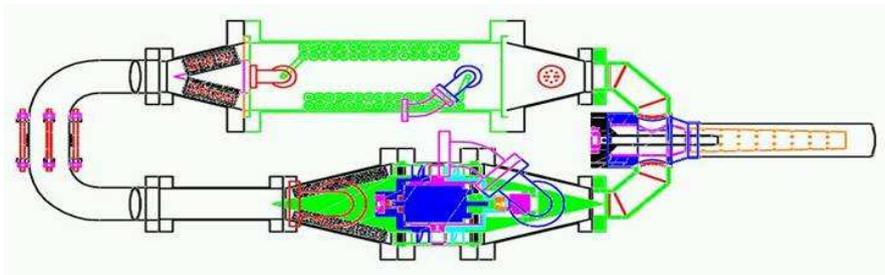
- along the focal planes of each of the toroid octants
- Back and front coincidence to eliminate neutrals
- 4 PMTs in low field region

Detect recoil protons with  $62^\circ < \theta < 78^\circ$



# G0 target

- 20 cm LH<sub>2</sub> cell, 250 W heat load from beam at 40  $\mu$ A
- High flow rate to minimize target density fluctuation
- magnet focal planes are independent on the point of interaction in the target



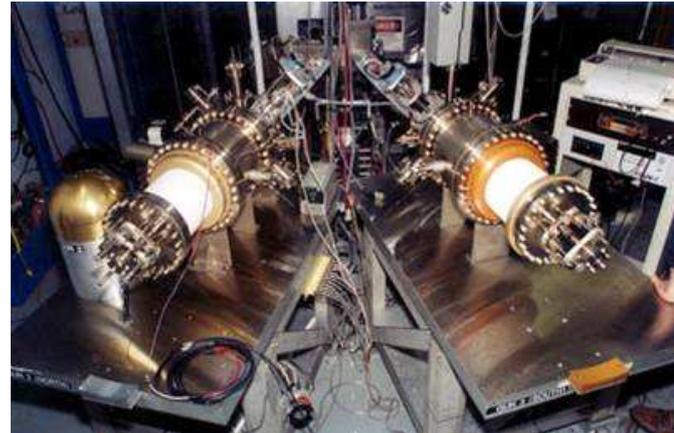
# G0 beam

## 32 ns between pulses - 31 MHz time structure

- 1/16 of usual CEBAF 499 MHz time structure
- higher charge per bunch
- required new Ti:Sapphire laser in polarized electron gun
- Helicity reversed with a pseudo random sequence: quartet structure + - -+ or - + +-
- Helicity held constant every 1/30 s (MPS)

## Beam quality achieved (IN-OUT):

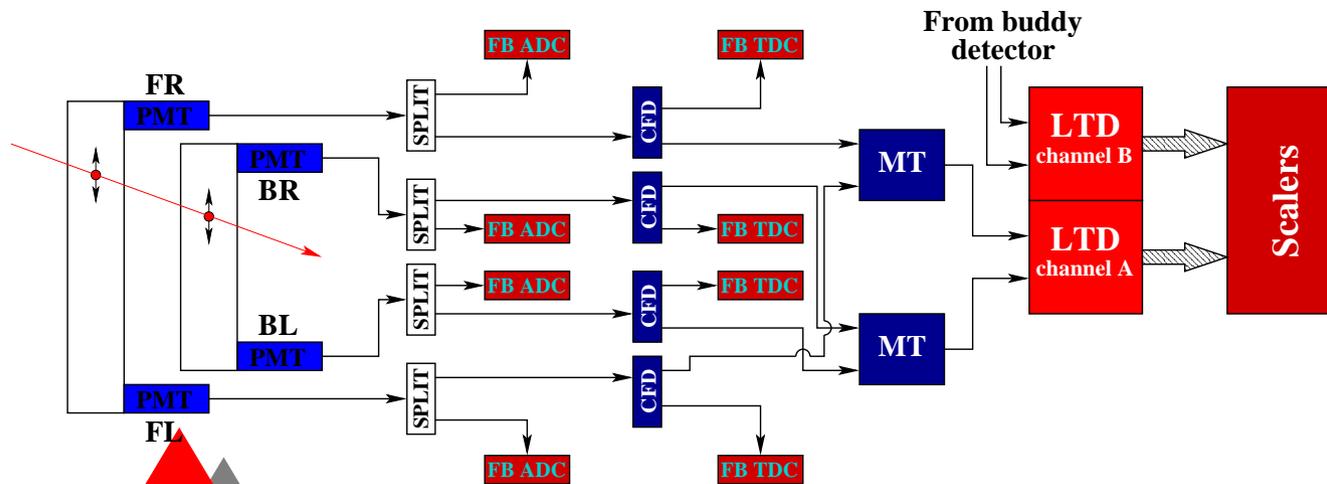
- $\Delta X = 6 \pm 4 \text{ nm}$
- $\Delta Y = 8 \pm 4 \text{ nm}$
- $\Delta x - \text{angle} = 2 \pm 0.3 \text{ nrad}$
- $\Delta y - \text{angle} = 3 \pm 0.5 \text{ nrad}$
- $\Delta I/I = -0.3 \pm 0.3 \text{ ppm}$
- Energy diff  $\Delta E = 58 \pm 4 \text{ eV}$



# G0 Electronics

Custom electronics design to provide high rate histogramming of time-of-flight spectra

- NA: meantimer → LTD → scalers - 1 ns resolution - 24 ns window
- FR: meantimer → DMCH - 0.25 ns resolution - 32 ns window
- data read out every MPS (captured between helicity periods)
- additional Fastbus data (ADC, TDC) at lower rate (3 FB events per MPS) for systematic studies (singles rates ...)



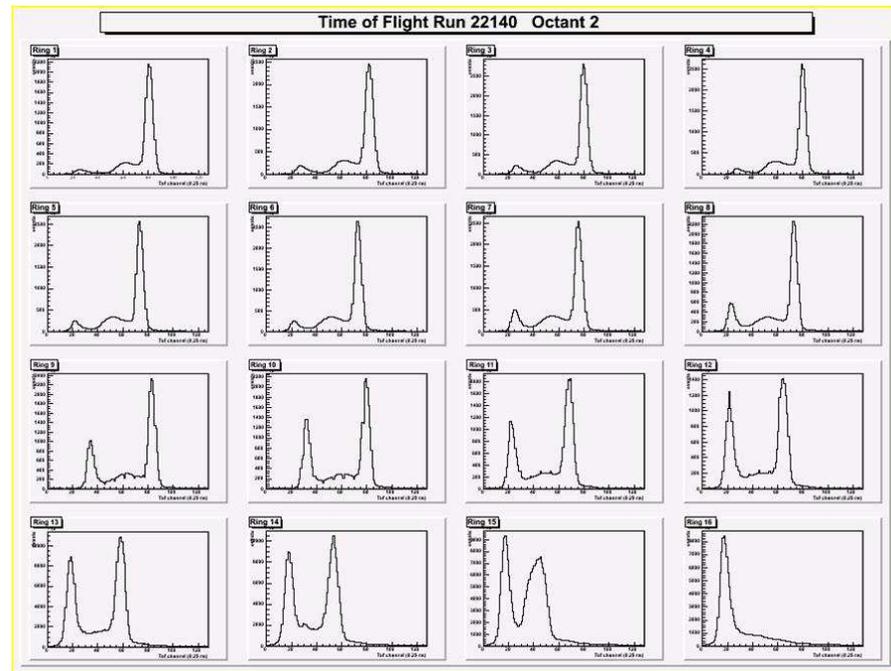
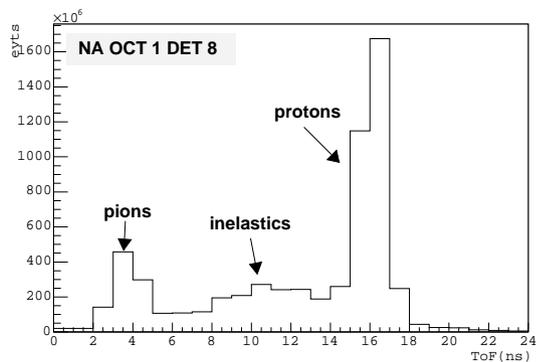
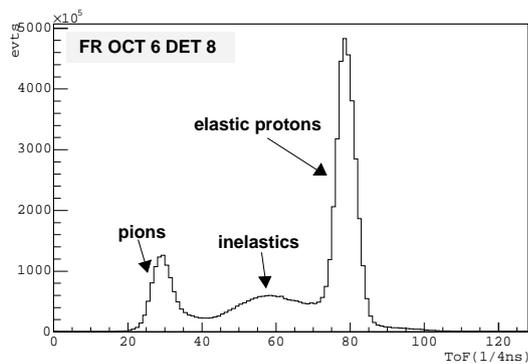
# The forward angle data

Detecting recoil protons

16 × 8 time-of-flight spectra

event rate of about 1MHz

PID by time-of-flight:  $p$  20 ns,  $\pi^+$  8 ns



# Analysis procedure

- For each helicity state (MPS)  $\rightarrow Y$  [ $\#counts/nC$ ]
- For each "good" quartet (e.g. + - - +)  $\rightarrow$   
$$A = \frac{Y_{1++}Y_{2+-} - Y_{1--}Y_{2-+}}{Y_{1++}Y_{2++} + Y_{1--}Y_{2--}}$$
- Weighted average of all quartets asymmetries

Not so easy....

- Deadtime corrections  $Y_c = \frac{Y_{raw}}{1-f}$ ,  $f = 10 - 15\%$
- Correct asymmetry for helicity correlated beam properties  $P_i$   
$$A_{corr} = A_{meas} - \sum_{i=1}^N \frac{1}{2Y} \left( \frac{\partial Y}{\partial P_i} \right) \Delta P_i \quad \Delta P_i = P_+ - P_-$$
- Leakage correction  $A'_{corr} = A_{corr} + \frac{I_L}{I_C} \left( 1 - \frac{Y_L}{Y_C} \right) A_{IL}$
- Background correction  $A_{sig} = \frac{A'_{corr} - A_{back} f_{back}}{f_{sig}}$
- Beam polarization/radiative corrections  $A_{phys} = \frac{A_{corr}}{P_{beam} R_{rad}}$

# Deadtime corrections

Electronic deadtime depending on rates  $\rightarrow$  helicity correlated

$$Y_m^\pm = (1 - f)Y_t^\pm$$

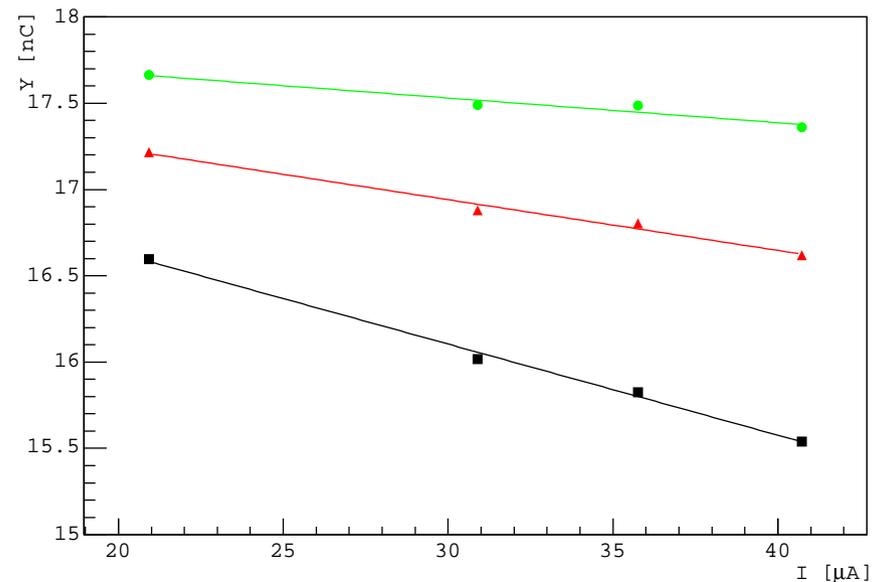
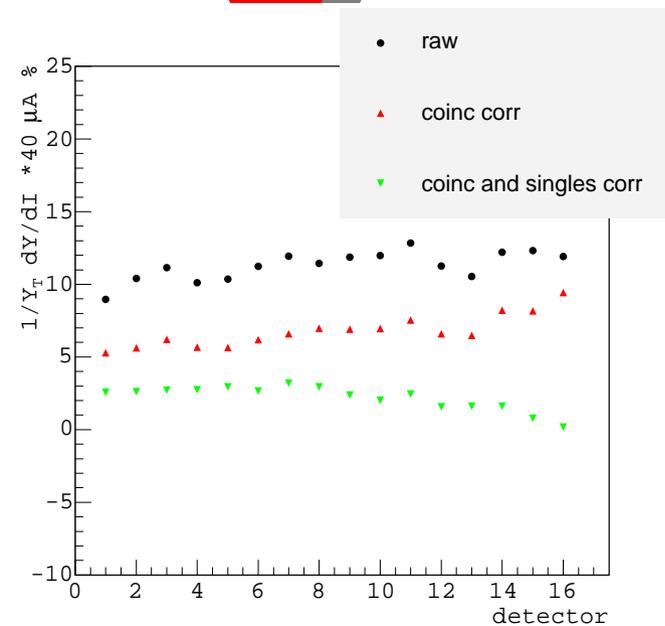
$$f \propto R$$

$f \sim 10\%$  before corrections

$$A_{false} \sim -f(A_Q + A_{phys})$$

Two "kinds" of deadtime:

- deadtime associated to coincidence rates  $\rightarrow$  NPN  $\rightarrow$  corrected for analytically
- deadtime due to "singles"  $\rightarrow$  need to measure/estimate singles rates  $\rightarrow$  FB data and/or scalers



# Leakage

Unanticipated effect: Leakage from Hall A,B lasers into Hall C

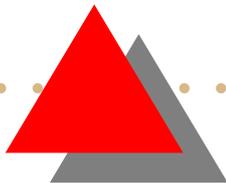
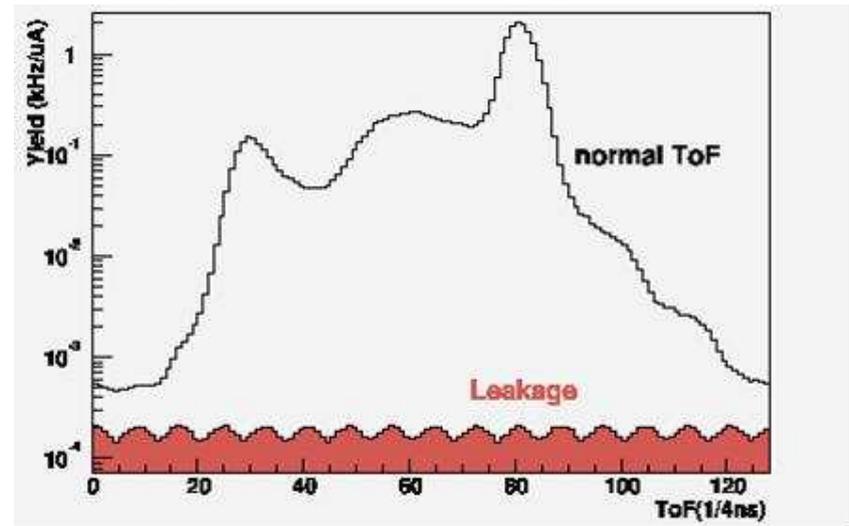
- Hall A,B beams and Hall C have different time structure (2 ns vs 32 ns)
- BCMs measure integrated (30Hz) A+B+C charge

$$A_{meas} = A_{phys} - \frac{I_L}{I_C} \left(1 - \frac{Y_L}{Y_C}\right) A_{IL}$$

$I_L, I_C$  leakage, Hall C currents

$A_{IL}$  leakage charge asymm

$Y_L, Y_C$  leakage, Hall C yields



# Leakage

To correct for this effect:

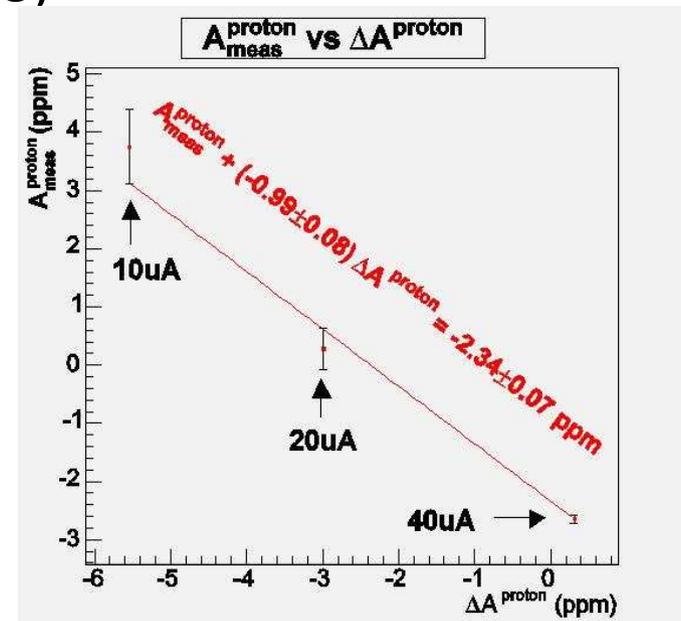
- measure the leakage asymmetry in a "forbidden" TOF region
- verify with direct leakage measurements by turning one laser on at the time (Lumi detectors)
- verify with low rates runs

Leakage current  $I_L = 40 \text{ nA}$

Charge asymmetry  $A_{IL} = 500 \text{ ppm}$

net systematic  $0.1 \text{ ppm}$

$$A_{phys} = A_{meas} + \Delta A$$

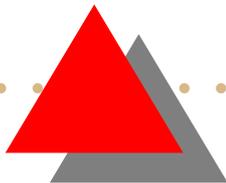
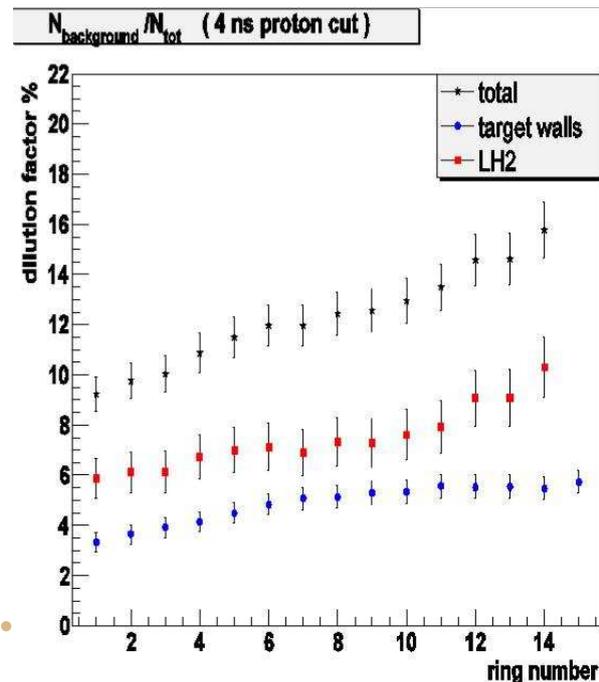
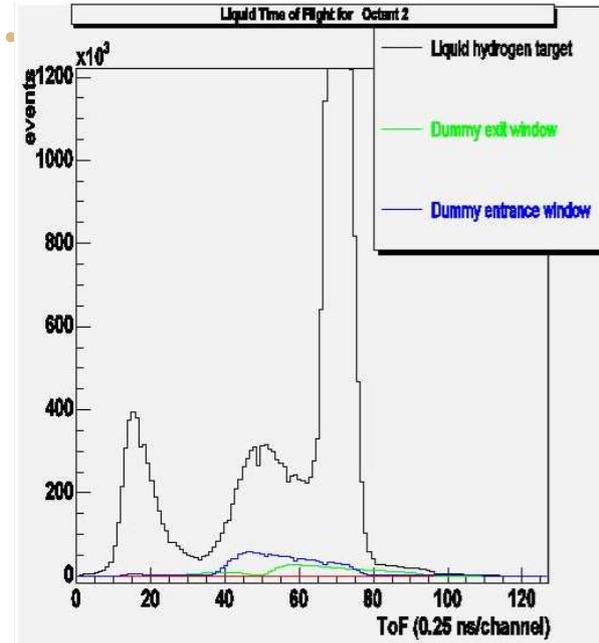


# Background correction

Need to measure dilution factor of elastic peak and the asymmetry of the background

- Data with full and empty targets, different pressures
- Data with dummy entrance and exit window
- Data with W radiator and dummy windows

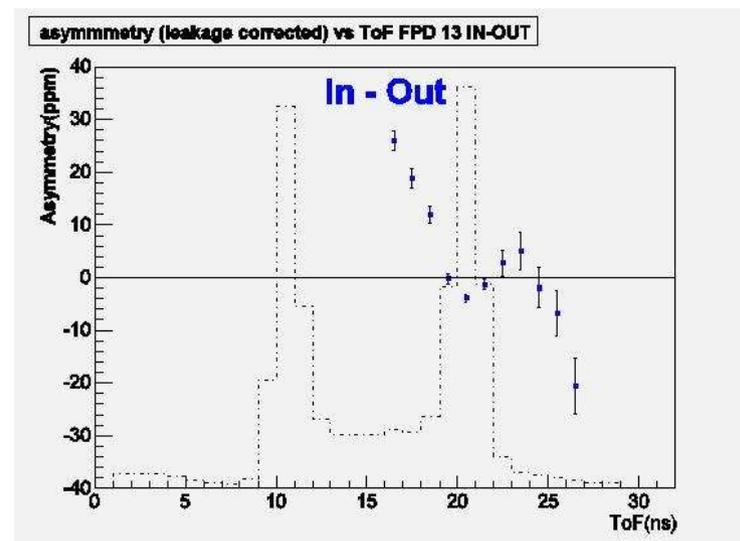
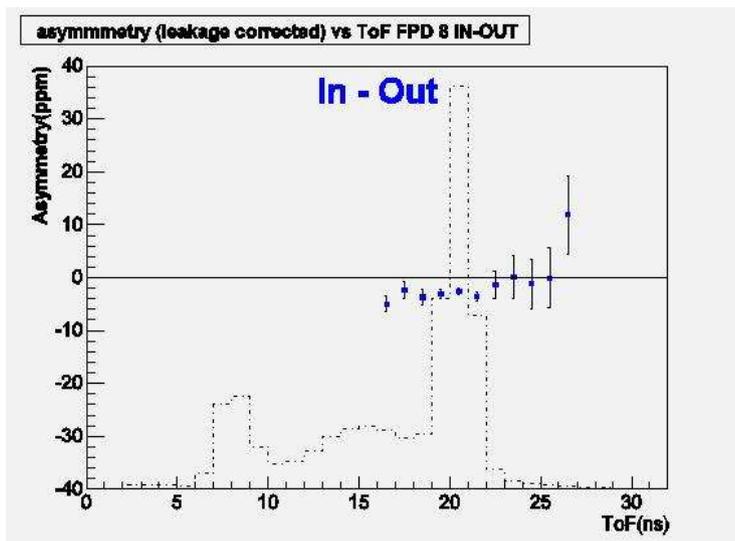
unfold background from target windows and the inelastic scattering from LH<sub>2</sub>



# Background correction

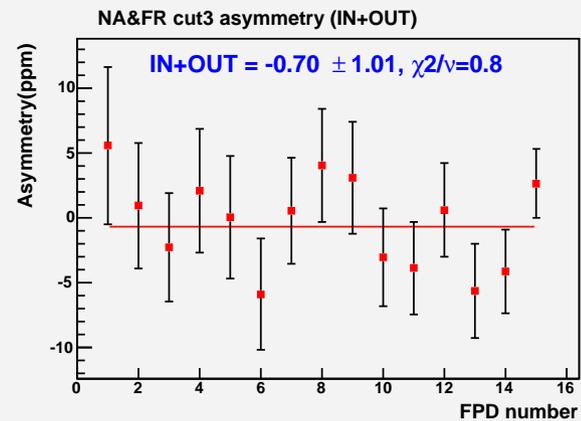
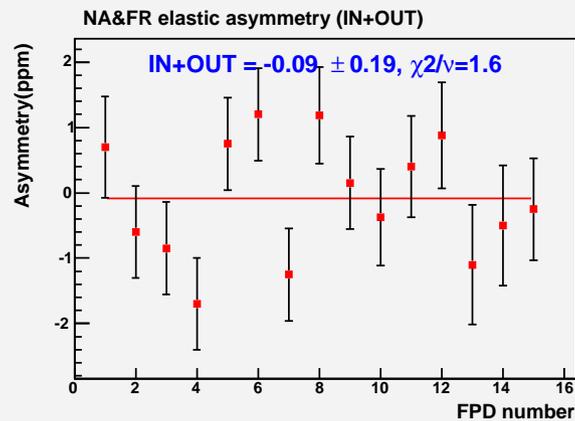
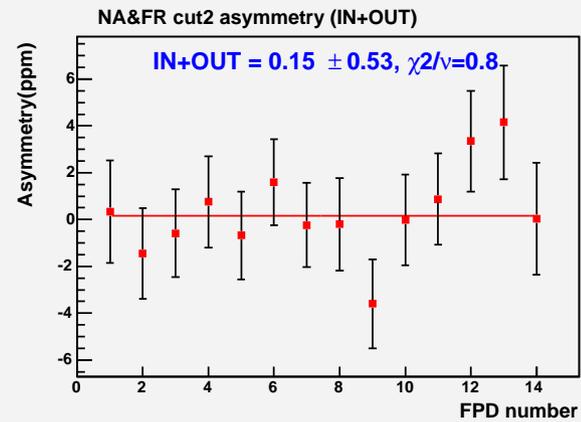
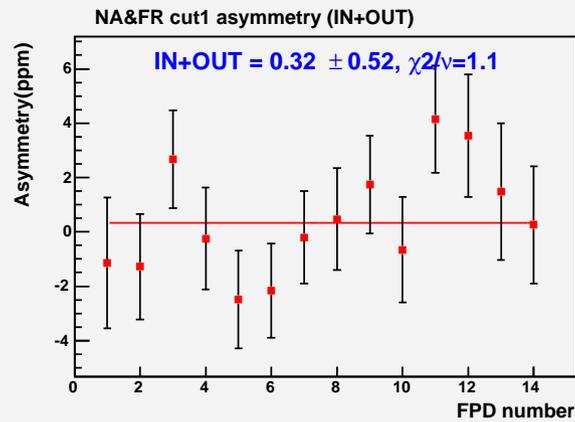
"Side-band" correction: Asymmetry and yield measured on either side of elastic peak  $\Rightarrow$  interpolation

- low/mid  $Q^2$  - error  $\sim$  2% - 5%
- high  $Q^2$  - present estimate 7 % - 20 %



# Results: IN+OUT

IN+OUT asymmetry of elastics and side-bands, 02/11-04/16



# Errors



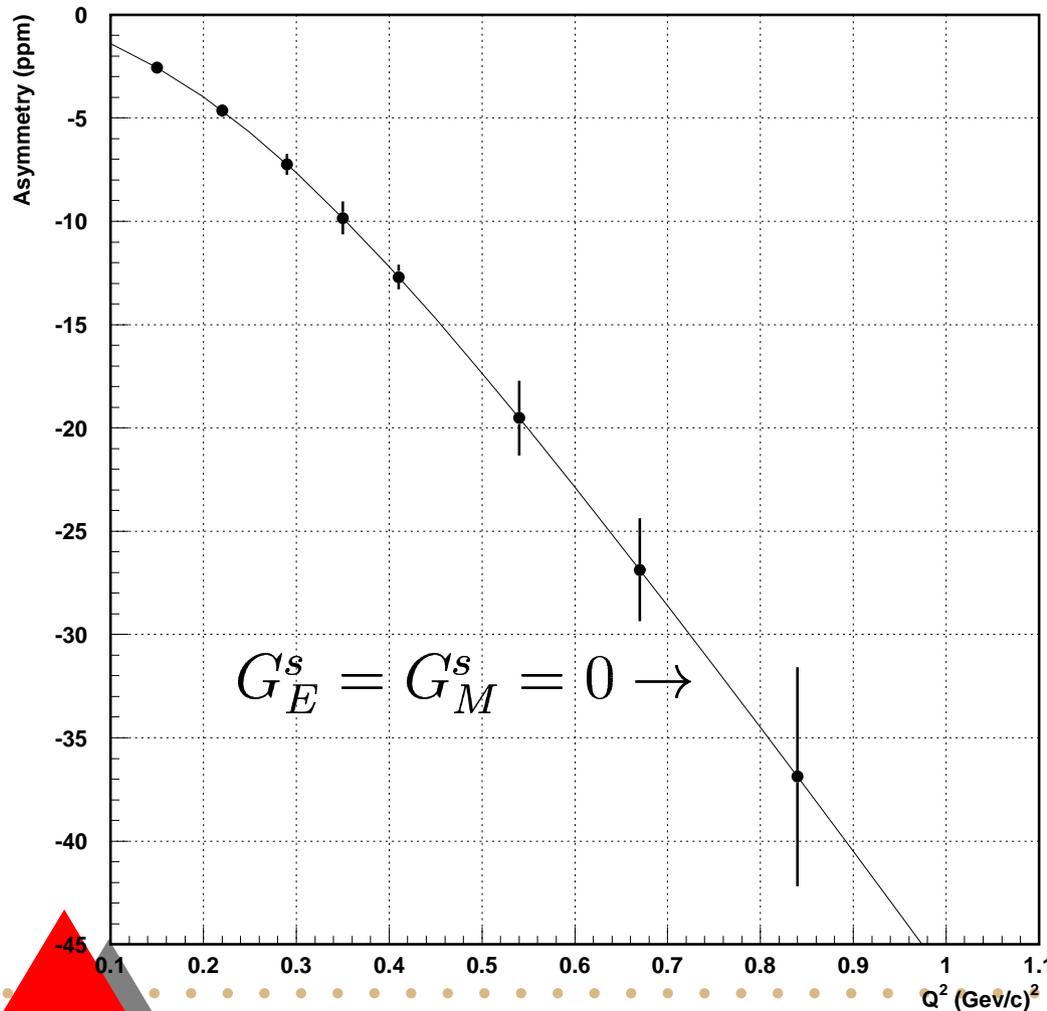
<i>source</i>	<i>error</i>
Deadtime correction	2 %
Beam param false asymm	0.01 ppm
leakage	0.10 ppm
Beam pol	2 %
Background	under study
Bin $Q^2$	1 %
Radiative corrections, EM form factors	to do

$G_E^s + \alpha G_M^s$  error probably dominated by background error, especially at high  $Q^2$



# Presently estimated uncertainties (statistical + systematic)

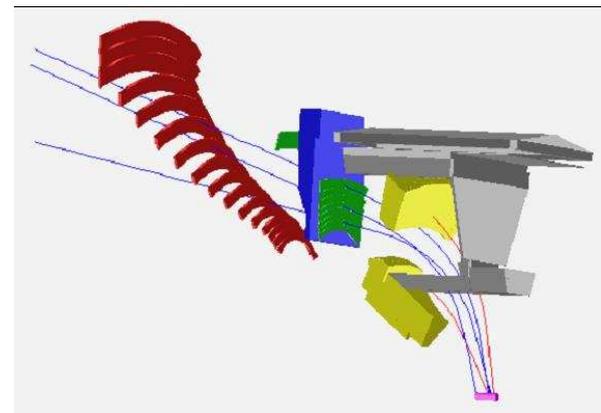
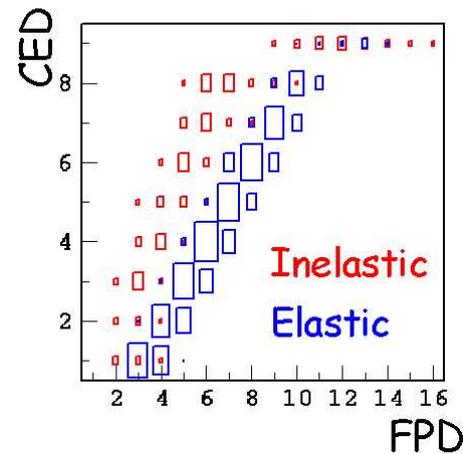
G0 expected precision (statistic + systematic)



# Future..

## Back-angle measurement to separate $G_E^s$ and $G_M^s$

- Detect scattered electron at  $\theta_e \sim 110^\circ$
- Back-angle  $Q^2$  have small variation in G0 acceptance
  - 3 energies 424, 576, 799 MeV for  $Q^2 = 0.3, 0.5, 0.8 \text{ GeV}^2/c^2$
- Need to extract  $G_A^e$ 
  - all measurements on both  $LH_2$  and  $LD_2$
- Setup turned around + new hardware and electronics
  - Cryostat Exit Detector (CED) to separate elastic and inelastic electrons
  - Cerenkov detector for pion rejection



First run - Fall 2005



# Conclusions

- G0 forward angle run complete
- Analysis well underway
- Background likely to dominate the errors
- Back angle run in Fall 2005
- In a few years we will have  $G_M^s$  and  $G_E^s$  for several  $Q^2$

